

Lessons from Pluto AMR

Adaptive mesh refinement in numerical astrophysics

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[with David Melon Fuksman, Mario Flock, Giancarlo Mattia, Andrea Mignone & gPLUTO devs]

AthenaK Summer School 2026

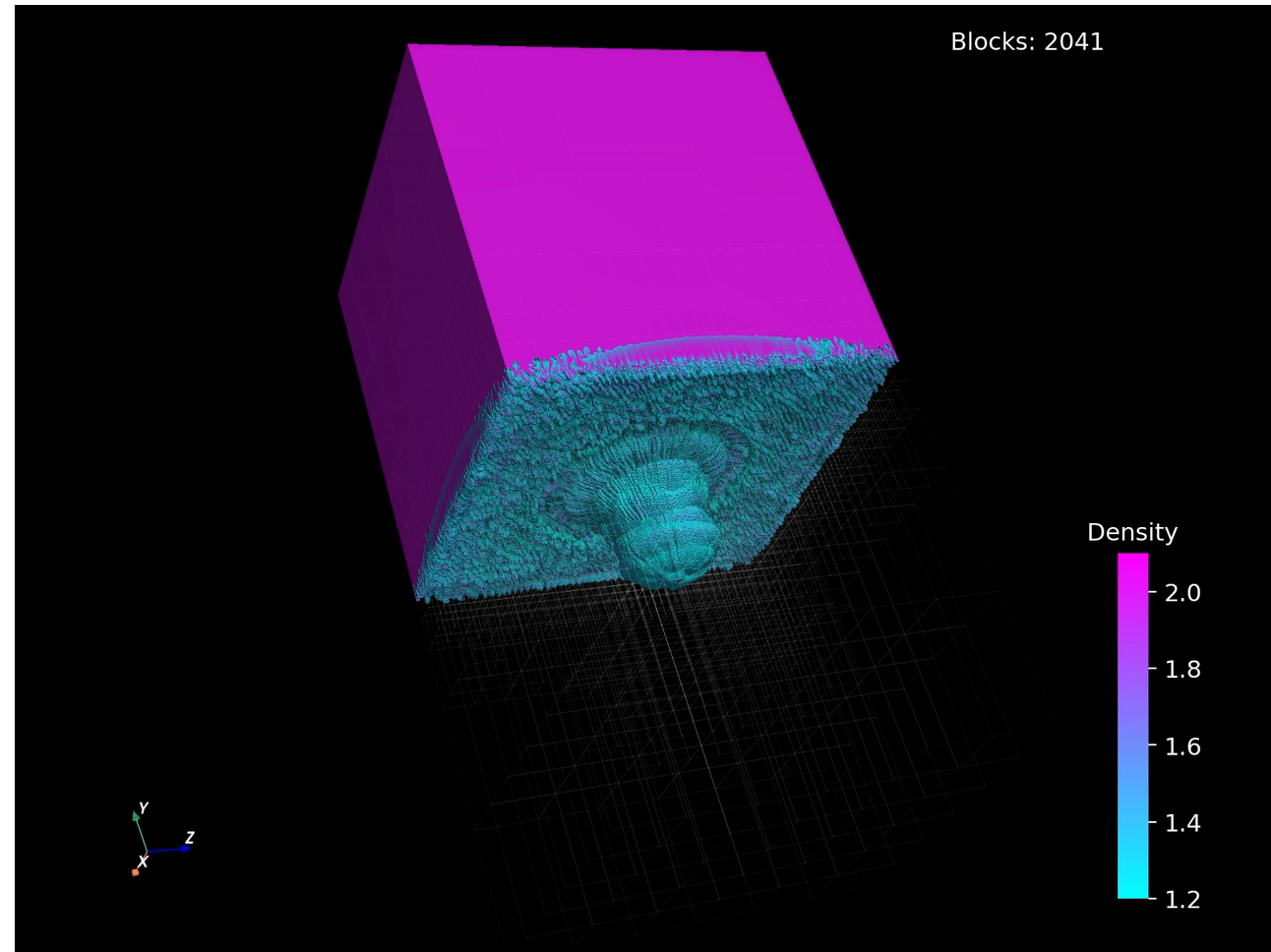
July 10, 2026



Motivation

Why adapt the mesh at all?

- You already know why! **See Jim's talk from Monday.**
- a uniform grid at that resolution is *hopelessly expensive* so put cells where the physics is!



3D Rayleigh-Taylor instability

[Melon Fuksman w gPLUTO]

Outline

Mesh refinement strategies

- Patch based
- Oct tree
- Moving meshes
- No mesh!

PLUTO → gPLUTO

same physics for GPUs is gBetter™ → oct-trees

[Melon Fuksman]

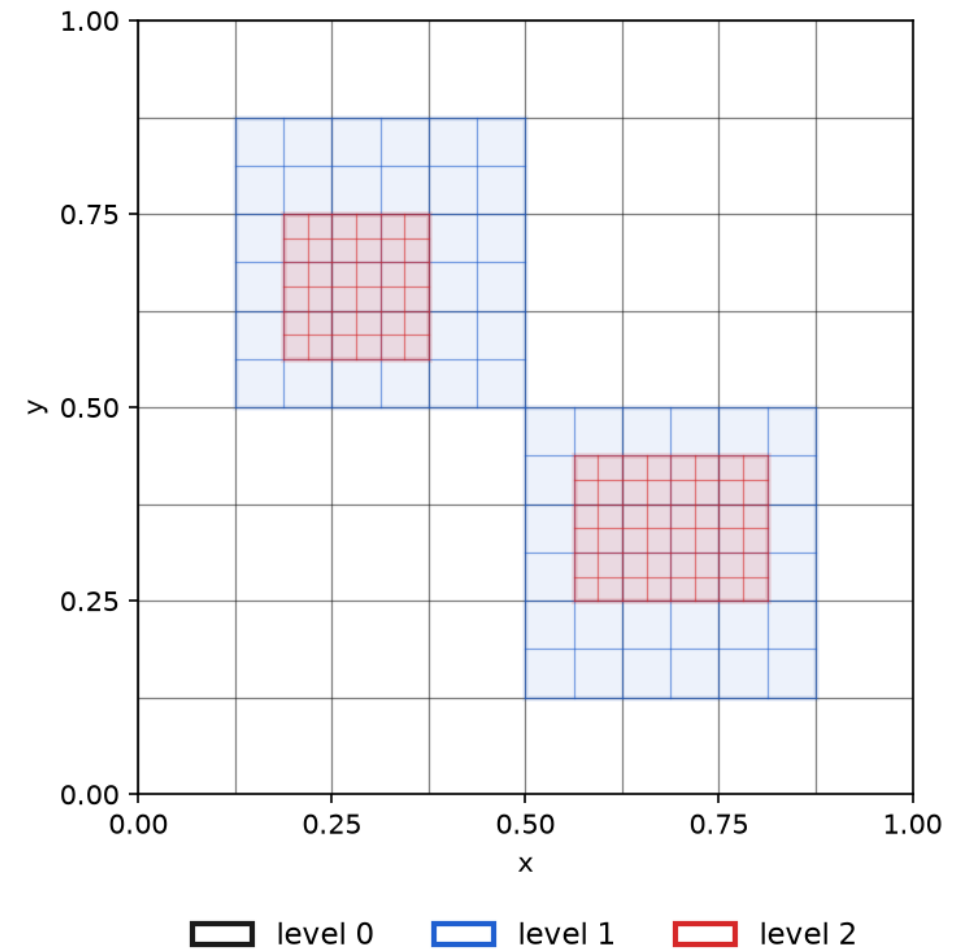
Staggered-grid AMR

$\nabla \cdot \mathbf{B} = 0$ for staggered grid



Patch based methods

- [Berger & Oliger 1984]
- whole domain is covered by a global coarse grid (level 0)
- refinement overlays rectangular patches on the regions that need resolution
- levels increase resolution by 2
- only restriction: a fine patch (level $l+1$) lies entirely within its parent (level l)
- **Trade-off**: keeping the boxes rectangular covers some “calm” space around the feature



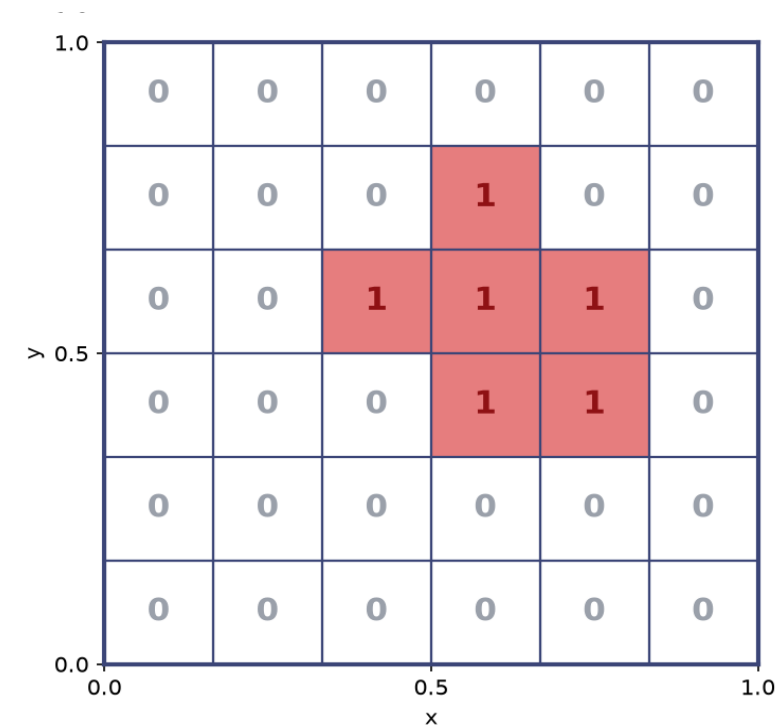
Refinement

- Richardson extrapolation [Berger & Oliger 1984]: criteria based on local estimation of the truncation error.

One step on the $2h$ grid (u_c) vs two steps on the h grid (u_f) estimates this:

$$\tau_{loc} \approx \frac{u_c - u_f}{2^p - 1} \quad (\text{order } p)$$

- **tag** cell 1 if $|\tau_{loc}| > \text{threshold}$, else 0
- Patches must cover the flagged cells $\rightarrow$ partitioning algorithms!



threshold $\varepsilon = 0.5$

Clustering algorithm

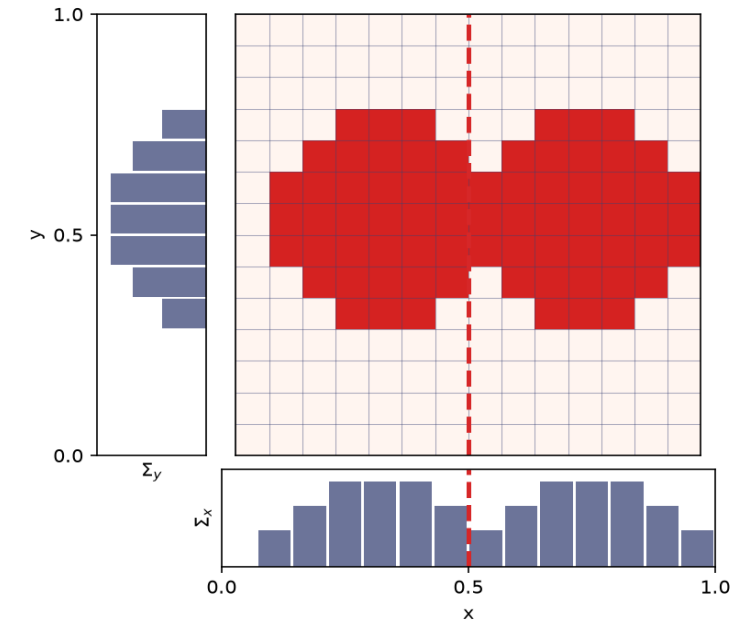
- [Berger & Rigoutsos 1991] turn the flagged cells into a few rectangles

- 1D signatures:

$$\Sigma_x(x) = \sum_y f(x, y), \quad \Sigma_y(y) = \sum_x f(x, y)$$

- **split** – at a **hole** ($\Sigma = 0$) or an **inflection point** (a zero-crossing of $\Delta^2 \Sigma$) to isolate clusters
- **stop** – recurse until every box meets the fill target:

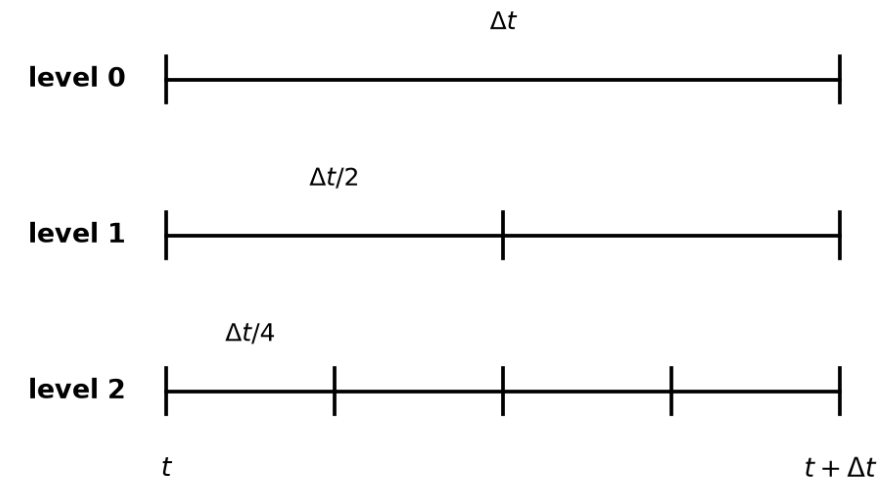
$$\eta = \frac{\sum f}{N_{\text{box}}} \geq \eta_{\text{target}}$$



split at a hole ($\Sigma_x = 0$) or at an inflection point ($\Delta^2 \Sigma_x$ zero-crossing)

Temporal sub-cycling

- Because of CFL constraint decouple the clocks
- **time BCs**: fine ghost zones between coarse time points are filled by **linear interpolation in space *and* time**



Refluxing

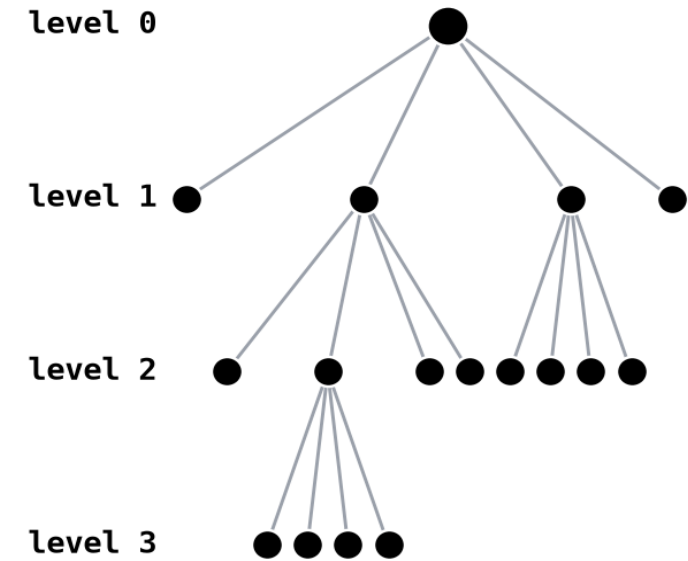
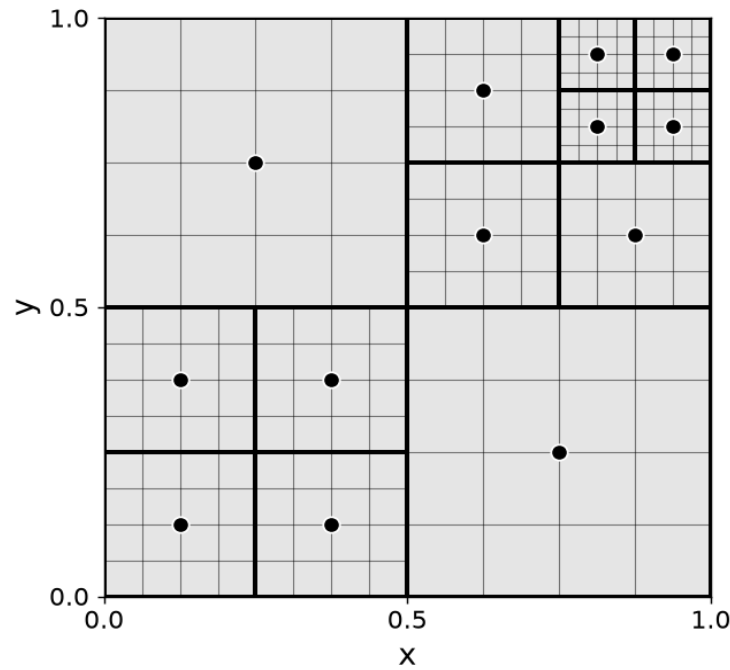
- [Berger & Colella 1989]
- coarse step and the r fine sub-steps compute *different* total fluxes across their shared face
- loss of conservation: uncorrected, mass & energy are created/destroyed at the patch perimeter
- Store the time-averaged fine flux minus the coarse flux:

$$\delta F_{i+1/2} = \frac{1}{r} \sum_{k=1}^r F_{\text{fine}}^{(k)} - F_{i+1/2}^C$$

- trust the fine grid: apply δF to the coarse cells just outside the patch → strict conservation

Tree-based AMR

- The fundamental unit of the grid is a contiguous block, the **MeshBlock**
- A block flagged for refinement splits into 2^d children → every child keeps the same number of points
- Blocks never overlap, so coarse and fine levels only need to communicate at their shared boundaries
- A logical tree representation → neighbours are easy to find



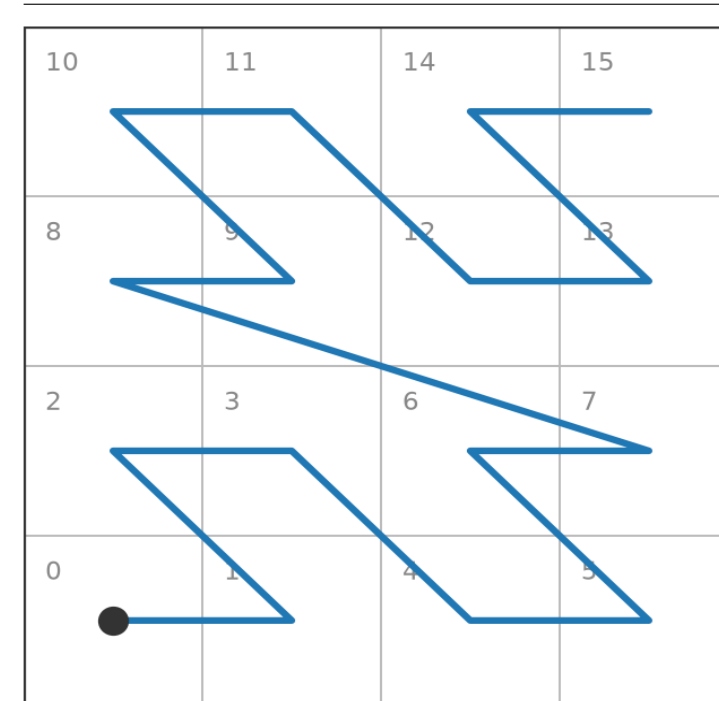
refinement

Space-filling curves and Morton order — Part I

In a large problem with thousands of meshblocks, these must be sensibly distributed across many GPUs. To keep inter-block communication efficient, neighbours should be assigned together whenever possible.

A *space-filling curve* maps 3D physical coordinates onto a 1D line while preserving spatial locality.

- **The Hilbert curve:** benchmark SFC: strictly continuous, folding on itself without ever crossing or jumping
- Cut the 1D curve into equal segments and every GPU gets the most compact volume possible → neighbors share face
- **Cons:** the Hilbert index needs stateful rotations and sequential bitwise XORs → regridding is expensive
- **Morton / Z-order:** just interleave the bits of (i, j, k) and level info
- cheap to compute
- occasional long “Z” jumps



Space-filling curves and Morton order — Part II

Shift each block onto the finest grid (by its level ℓ), then interleave bits of (i, j, k) into one **Morton key**:

$$\text{key} = \text{morton}(c \ll (L - \ell))$$

leaf blocks (i, j) with level:

[$(0,0)L_0$, $(1,0)L_0$, $(0,1)L_0$, $(1,1)L_0$, $(2,0)L_0$, $(4,2)L_1$, $(4,3)L_1$, $(5,3)L_1$,
 $(10,4)L_2$, $(11,4)L_2$, $(10,5)L_2$, $(11,5)L_2$]

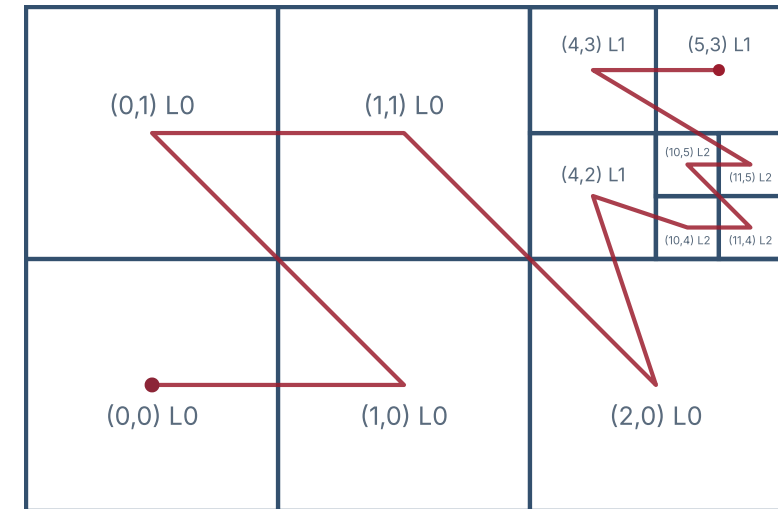
Morton keys

[0, 16, 32, 48, 64, 96, 104, 108, 100, 101, 102, 103]

sort!

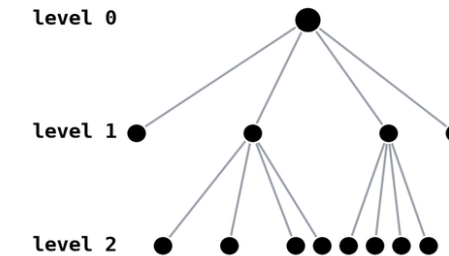
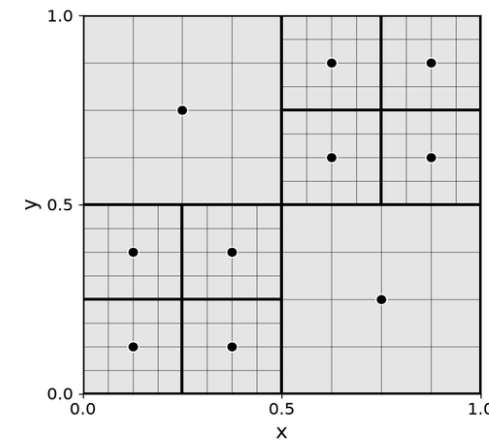
[0, 16, 32, 48, 64, 96, 100, 101, 102, 103, 104, 108]

the sorted keys **are** the Z-order – the flat linear octree



Neighbor finding

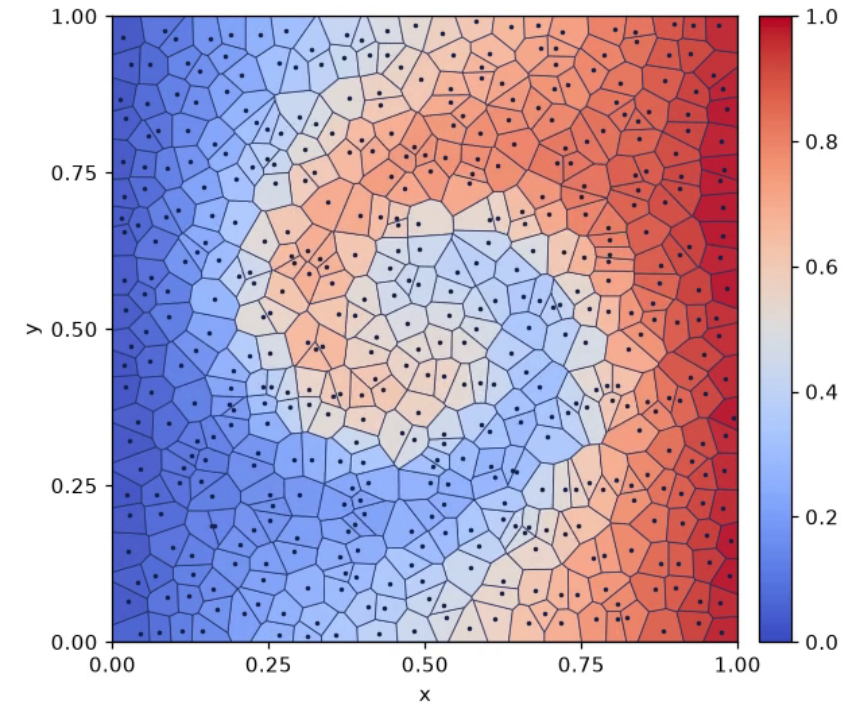
- Ghost exchange, flux matching and the 2:1 ripple all ask “**who is next to me?**”
- **Tree walking** – the pointer-tree way:
 - **walk up** – follow *parent* pointers to the smallest ancestor that also contains the neighbour
 - **walk down** – descend its *child* pointers toward the neighbour until you reach a leaf
 - $O(\log N)$ pointer chases per query, scattered in memory → cache-unfriendly
- **Flat oct-tree** – no traversal! step the coords $(i, j, k) \rightarrow (i-1, j, k)$, recompute the **level-aware Morton key**, and probe the hash table → the neighbour’s address in $O(1)$
- **across levels** – if the key’s **parent** is a leaf → the neighbour is **coarser**; if finer keys fill its **span** → **finer** (2:1 balance allows only these)



Moving meshes

A different philosophy: stop refining a *fixed* grid
– let the mesh move with the fluid.

- middle ground between fixed grid and mesh-less methods → has Galilean invariance and can handle shocks better than SPH.
- cells = **Voronoi** regions around points that advect at the flow velocity [Springel 2010]
- the tessellation rebuilds every step: topology changes for free
- resolution **follows the mass** “natural refinement

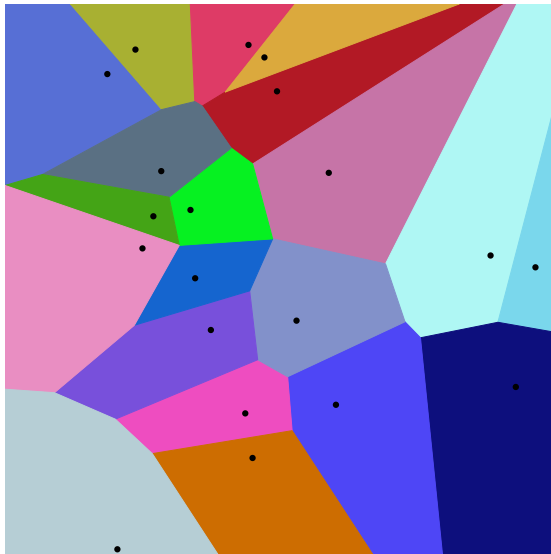


the mesh advecting with the flow

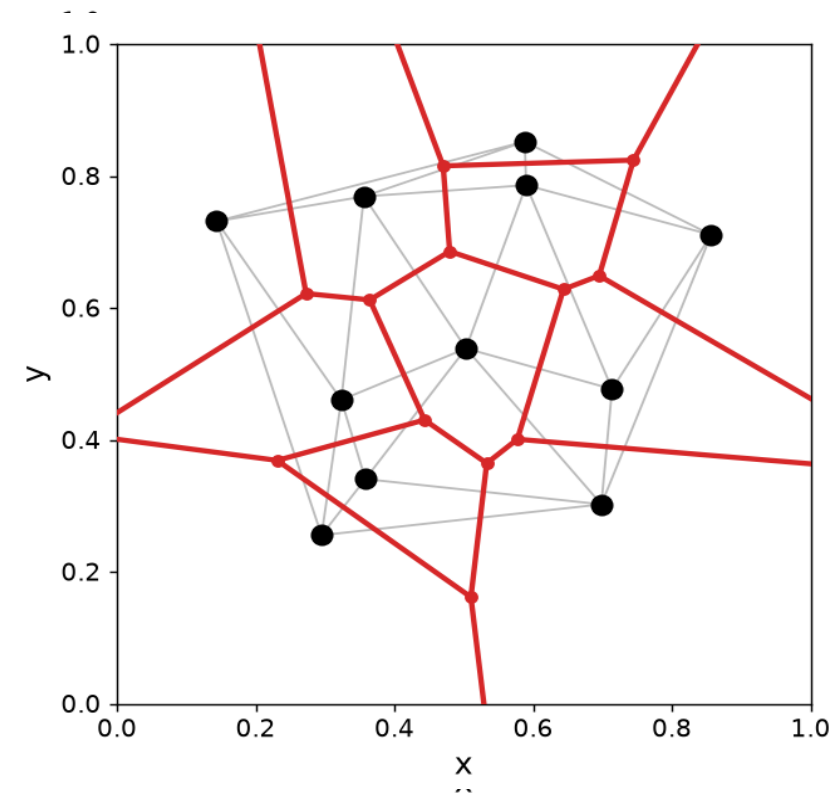
Building a Voronoi mesh

A Voronoi mesh tiles space into cells – each cell is everything closer to its generator r_i than to any other point.

In practice you build its dual, the **Delaunay triangulation**.



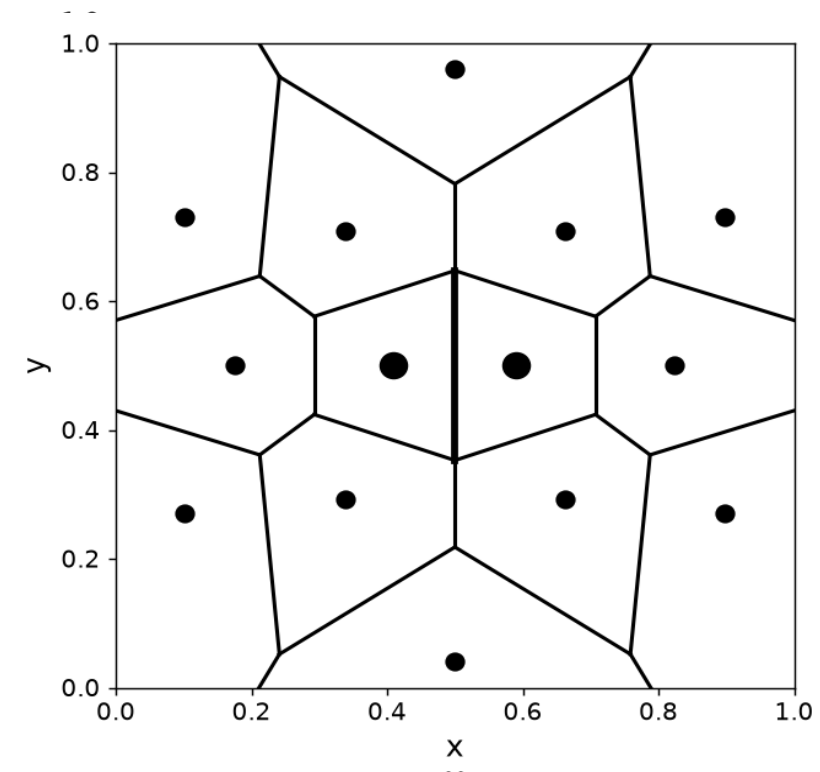
Voronoi diagram, Wikimedia Commons



Delaunay triangulation and its dual

Refinement and derefinement — I

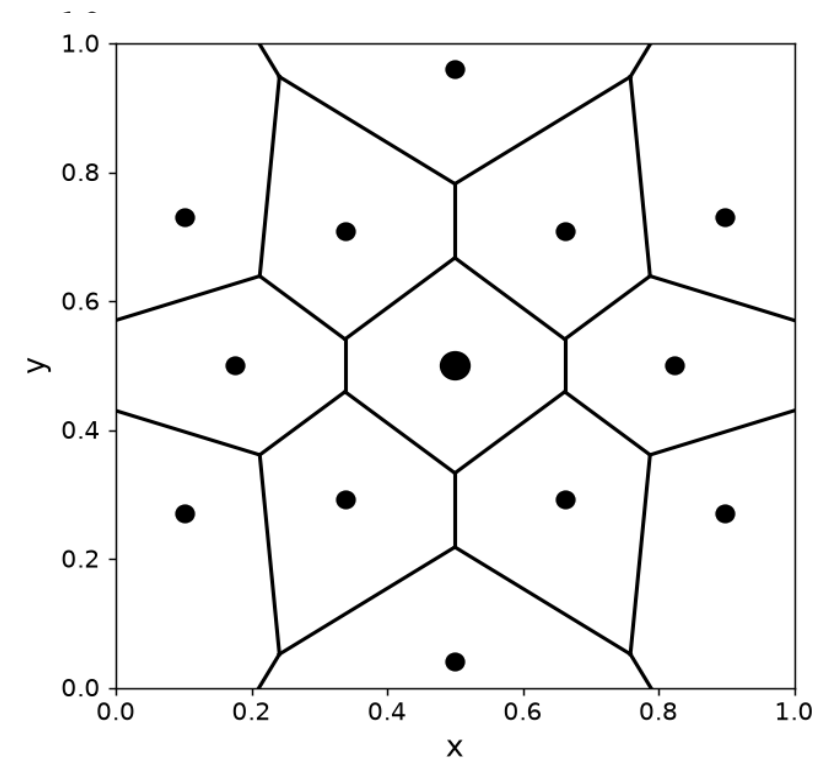
- Keep cell masses approximately equal
- Check against a threshold mass
- **Point splitting:** delete the flagged parent, inject **two** new points in its place, displaced slightly from the centre along the cell's longest dimension
- **Conserve the physics:** mass, momentum and internal energy split in half
 $m_{\text{child}} = \frac{1}{2} m_{\text{parent}}, \quad \mathbf{p}_{\text{child}} = \frac{1}{2} \mathbf{p}_{\text{parent}}$



one generator $\rightarrow$ two points along the long axis $\rightarrow$ the cell splits

Refinement and derefinement — II

- Just exact reverse
- **Point merging:** delete the two flagged neighbouring points and inject one point at their centre of mass
- Sum the children back into the parent



two points → merge to one at the centre of mass → the cells fuse

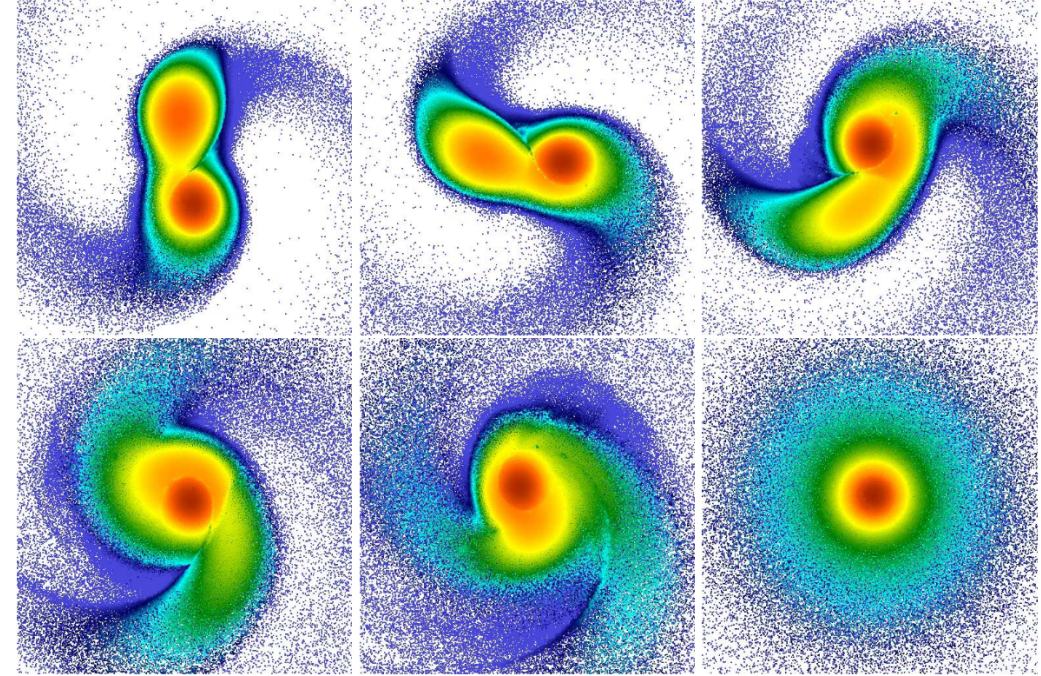
Disadvantages — expensive Voronoi rebuild each step · memory-intensive (connectivity + geometry) · mesh noise from irregular, moving cells

Meshless/SPH approaches

- No “mesh refinement” as there is no mesh! The particles are the fluid.
- Galilean invariance is perfect; mass, momentum conserved by construction.
- Empty space costs zero memory and compute.

Some cons:

- low-order & noisy; artificial viscosity for shocks
- $\nabla \cdot \mathbf{B}$ by cleaning, not CT



[Diehl et al. 2008] – $q = 1.3$ DD (double-degenerate) merger
with **NuGrid**

The SPH kernel

- Each particle carries fixed mass m_i and an interaction radius or the “smoothing length” h .
- Particles move with the local fluid velocity.

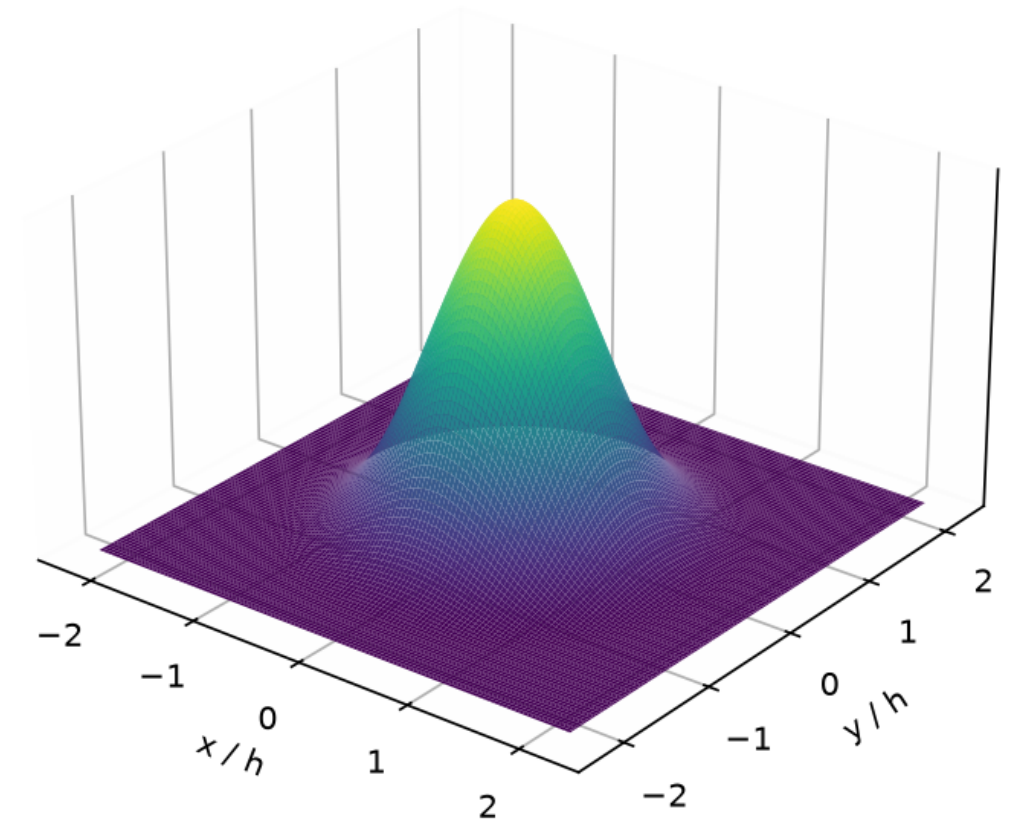
A quantity is a **kernel-weighted average** [Gingold & Monaghan 1977, Lucy 1977]:

$$\langle A(\mathbf{r}) \rangle = \int A(\mathbf{r}') W(|\mathbf{r} - \mathbf{r}'|, h) d\mathbf{r}' \longrightarrow \langle A \rangle_i = \sum_j m_j \frac{A_j}{\rho_j} W_{ij}$$

The equations of evolution are functions of W and its derivatives ∇W [Rosswog 2009, 2015]:

$$\rho_i = \sum_j m_j W_{ij}, \quad \frac{d\mathbf{v}_i}{dt} = - \sum_j m_j \left(\frac{P_i}{\rho_i^2} + \frac{P_j}{\rho_j^2} \right) \nabla_i W_{ij}$$

$$\frac{du_i}{dt} = \sum_j m_j \frac{P_i}{\rho_i^2} (\mathbf{v}_i - \mathbf{v}_j) \cdot \nabla_i W_{ij} \quad W_{ij} \equiv W(|\mathbf{r}_i - \mathbf{r}_j|, h)$$

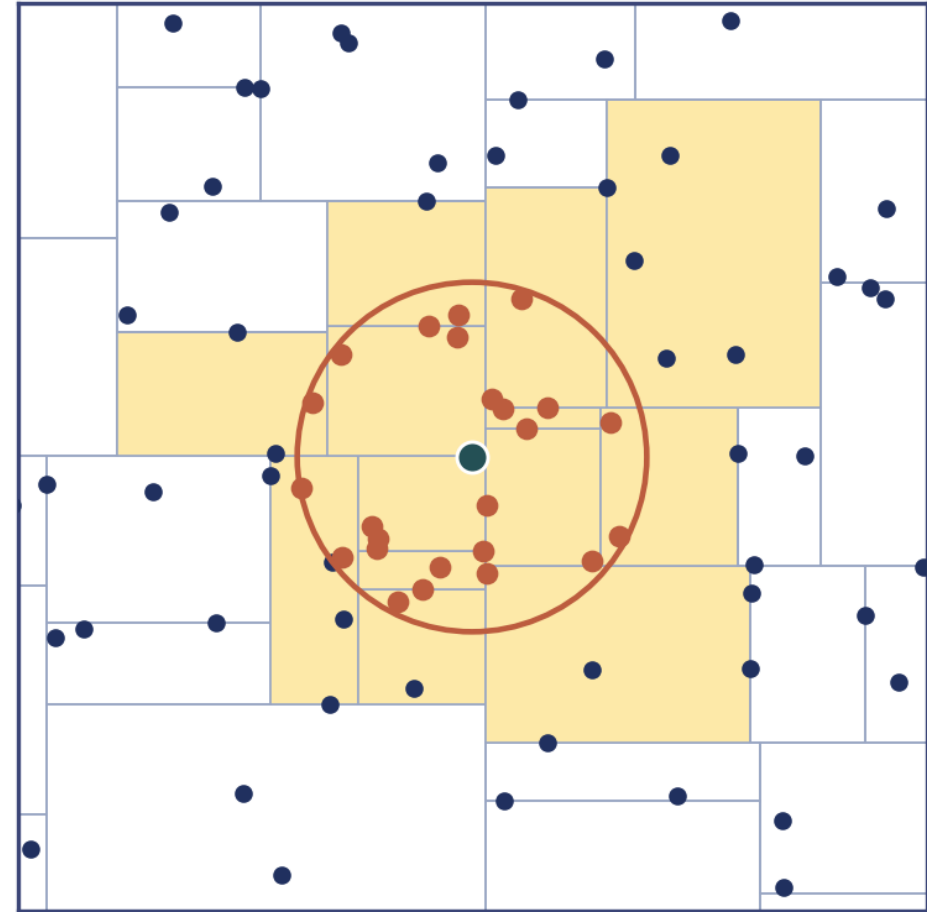


A cubic-spline kernel $W(r, h)$

Finding neighbours

Every particle needs the $\sim N_{\text{SPH}}$ others inside its kernel support $\lesssim 2h$, but

- distance calculation $O(N^2)$ expensive!
- sort the particles into a spatial tree; walk only the cells overlapping the $2h$ sphere $\rightarrow O(N \log N)$
- Rosswog et. al. uses a **recursive coordinate bisection (RCB) tree** [Gafton & Rosswog 2011]



Searches only within cells overlapping the $2h$ circle

A comparison

	Patch based (AMReX, Chombo)	Block oct-tree (AthenaK, gPLUTO)	Moving mesh (AREPO)	SPH (SPHINCS_BSSN)
refinement unit	arbitrary rectangle	fixed block “meshblock”/“amrblock”	Voronoi cell	particle
adaptivity	discrete, flagged	discrete, per block	continuous, co-moving	continuous, adaptive h
data structure	patches + clustering	oct-tree of blocks	Voronoi tessellation	neighbour lists
Galilean-invariant	X	X	✓	✓
$\nabla \cdot \mathbf{B} = 0$	CT + reflux	staggered CT + reflux	cleaning	cleaning

The PLUTO code

Solves a system of conservation laws [Mignone+ 2007]:

$$\frac{\partial \mathbf{U}}{\partial t} + \nabla \cdot \mathbf{T}(\mathbf{U}) = \mathbf{S}(\mathbf{U})$$

For MHD, the conserved state and flux:

$$\mathbf{U} = [\rho, \rho \mathbf{v}, \mathbf{B}, E]^\top$$
$$\mathbf{T} = \begin{bmatrix} \rho \mathbf{v} \\ \rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + p_t \mathbf{I} \\ \mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v} \\ (E + p_t) \mathbf{v} - (\mathbf{v} \cdot \mathbf{B}) \mathbf{B} \end{bmatrix}$$

- PLUTO-AMR [Mignone+ 2012] – AMR via Chombo



PLUTO/gPLUTO devs and users @ MPIA, 2026

Public repo – gitlab.com/PLUTO-code/gPLUTO

Webpage – plutocode.ph.unito.it

PLUTO → gPLUTO

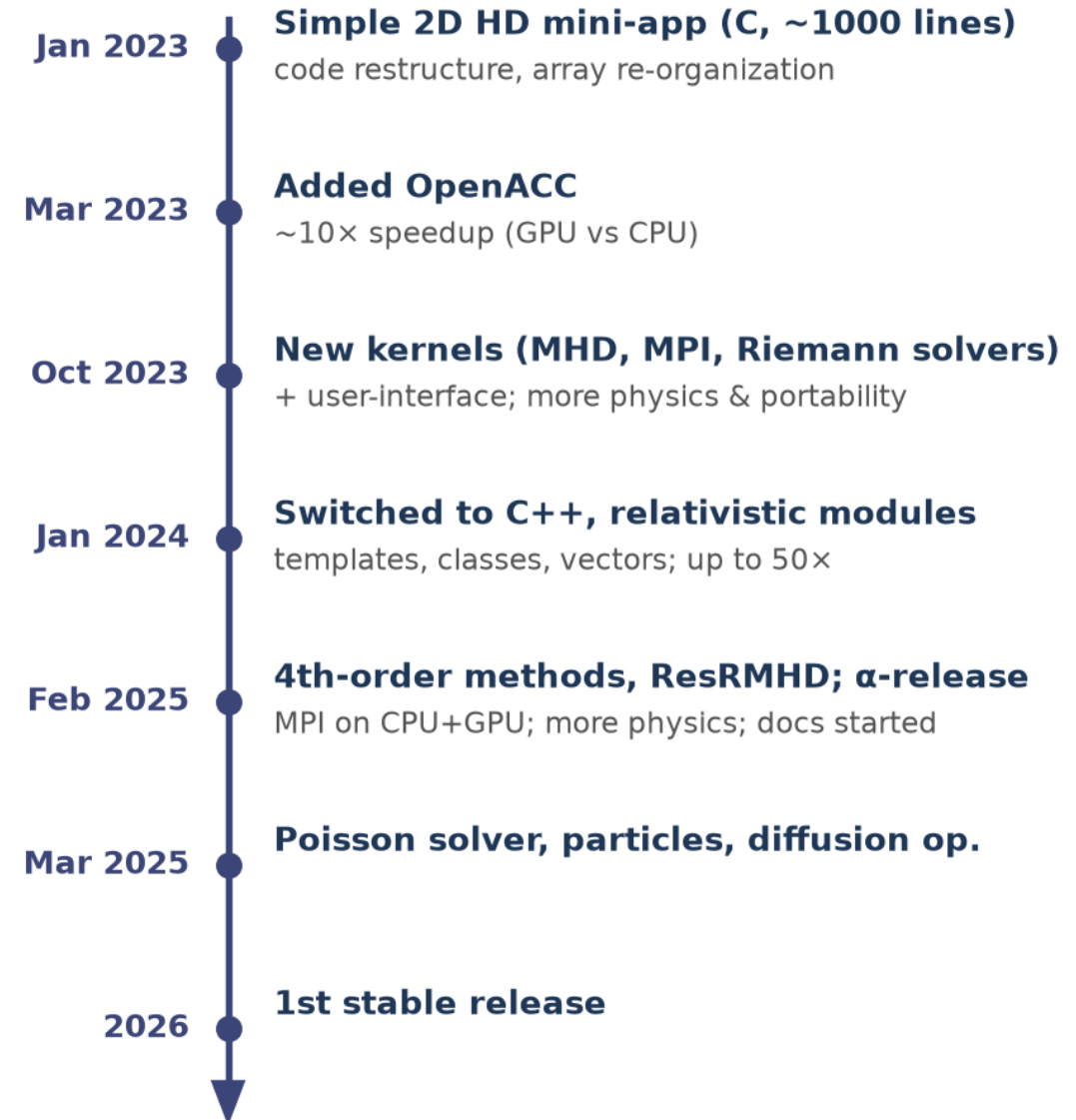
- Built within the EuroHPC *Scalable Parallel Astrophysical Codes for Exascale* (SPACE)
- Ready PLUTO for exascale on heterogeneous CPU + GPU



Directive-based port is easy and incremental

- one C++ source; annotate the hot loops for target hardware:

```
1 // OpenACC
2 #pragma acc parallel loop
3 for (i = 0; i < n; i++) update(i);
4
5 // OpenMP
6 #pragma omp target teams distribute parallel for
7 for (i = 0; i < n; i++) update_something(i);
8
9 // hierarchical: gang over blocks, vector within
10 #pragma acc parallel loop gang
11 for (b = 0; b < nblok; b++)
12 #pragma acc loop vector
13 for (i = 0; i < bs; i++) update_something(b, i);
```



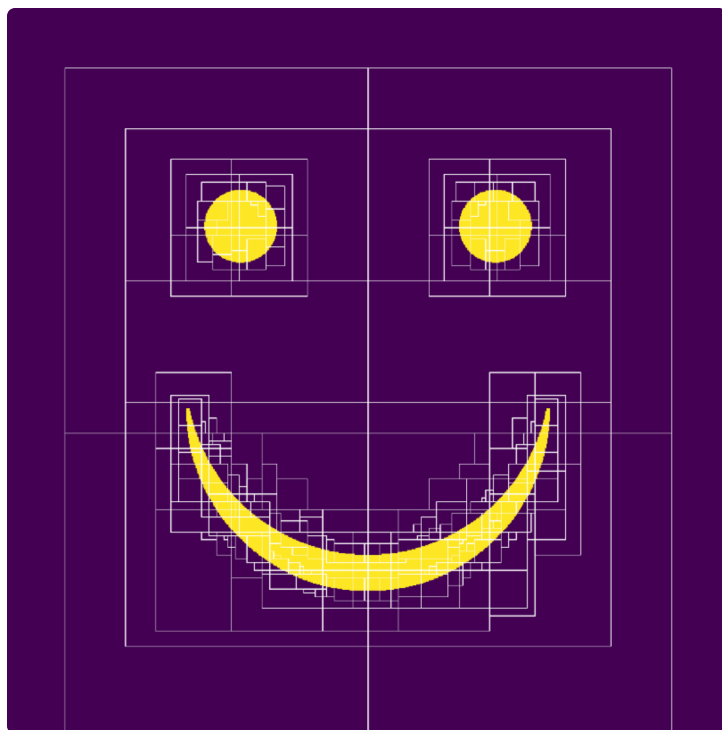
[Andrea Mignone]

Status of Port

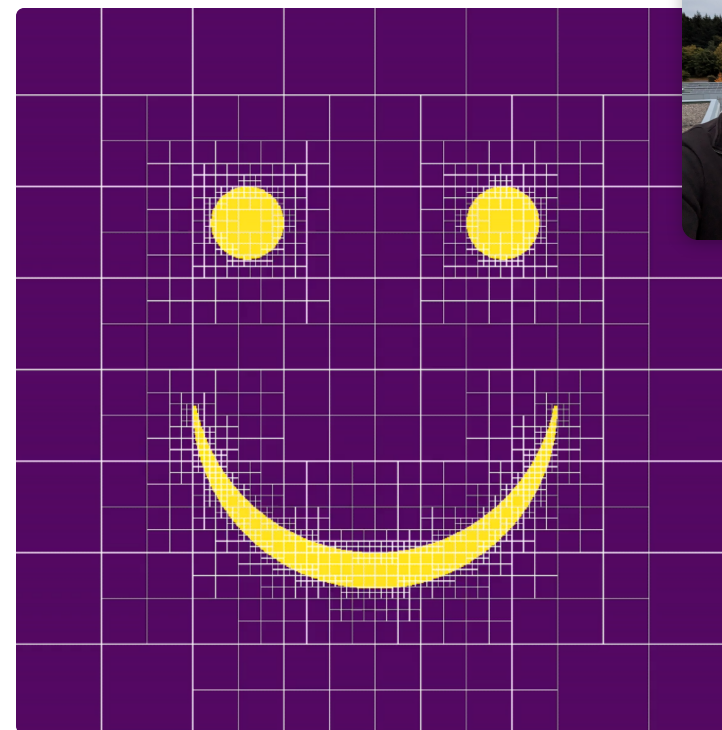
	PLUTO	gPLUTO	Pub		PLUTO	gPLUTO	Pub
Basic physics				Geometry			
HD (Euler)	✓	✓	✓	Cartesian	✓	✓	✓
MHD	✓	✓	✓	Cylindrical	✓	✓	✓
RHD (rel. hydro)	✓	✓	✓	Spherical	✓	✓	✓
RMHD (ideal)	✓	✓	✓	Cooling			
ResRMHD (non-ideal)	✓	✓	✓	Tabulated	✓	✓	✓
GRMHD	×	~ dev	–	Power law	✓	×	planned
Dissipation				SNEq	✓	✓	✓
Viscosity	✓	×	planned	MINEq	✓	×	–
Thermal conduction	✓	×	planned	Radiation / Particles			
Resistivity	✓	✓	✓	Radiation (classical)	✓	✓	✓
Ambipolar diffusion	×	~ wip	–	Radiation (relativistic)	✓	✓	✓
Hall MHD	✓	×	–	Cosmic rays	✓	~ wip	×
EoS				Dust	✓	×	planned
Ideal / isothermal	✓	✓	✓	Lagrangian particles	✓	✓	×
Synge gas	✓	✓	✓				
Non-constant γ	✓	×	–				

[A. Mignone]

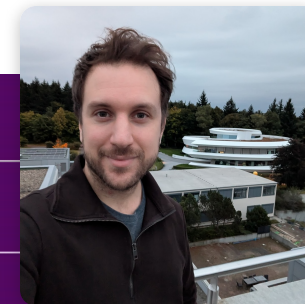
A new AMR for gPLUTO



Legacy PLUTO uses CHOMBO



gPLUTO oct-tree [David Melon Fuksman]



The AMR structure of gPLUTO

- gPLUTO implements the **oct-tree**! → this is easier than a patch-based approach!
- AMR support for **non-Cartesian coordinates** (cylindrical, spherical)
- **AMRBlock** – the basic unit of the grid, equivalent to AthenaK's **MeshBlock**
- **AMRRoot** – the class holding all AMRBlocks, arrays and methods
- Each MPI ranks gets a collection of **AMRBlocks** following the Morton curve
- Send and receive data through buffers happens through async MPI communications
- the tree structure is stored on **rank 0** for now – *changing soon*

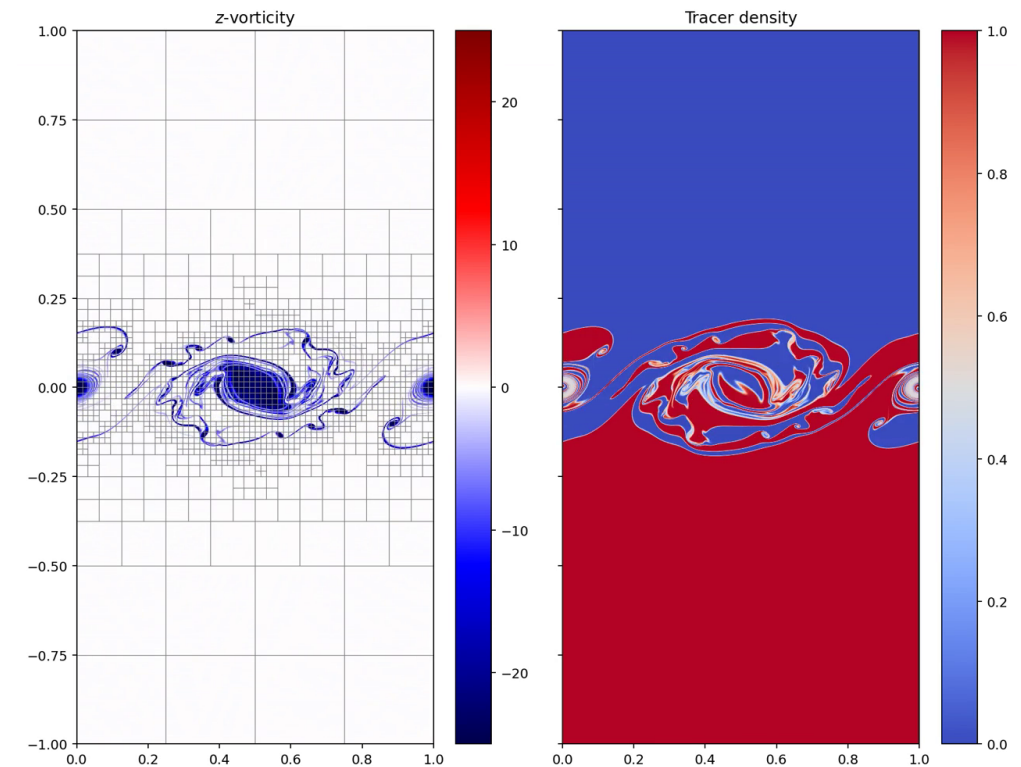
Three things happen each step

- 1 Fill boundaries**
prolong / restrict / copy
- 2 Integrate equations**
in each block
- 3 Ensure conservation**
refluxing

Refluxing at work — Kelvin-Helmholtz

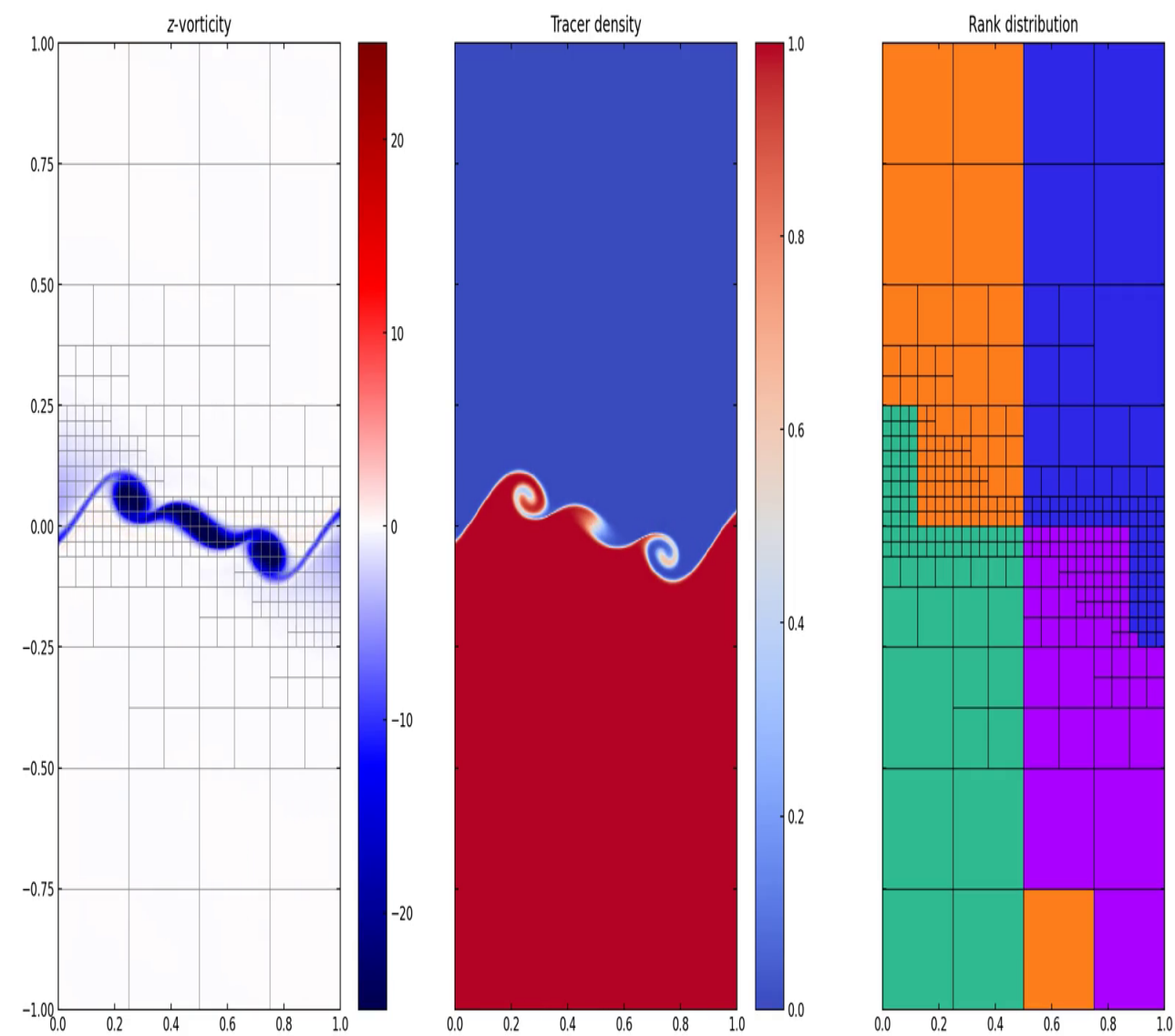
- the KH roll-up straddles coarse-fine boundaries, so it is a sharp test of **flux conservation**
- **fixed refluxing** keeps the vortices sharp and free of spurious features at the level jumps

[D. Melon Fuksman · gPLUTO]



Kelvin-Helmholtz, max level 5, Roe, fixed refluxing

Load balancing

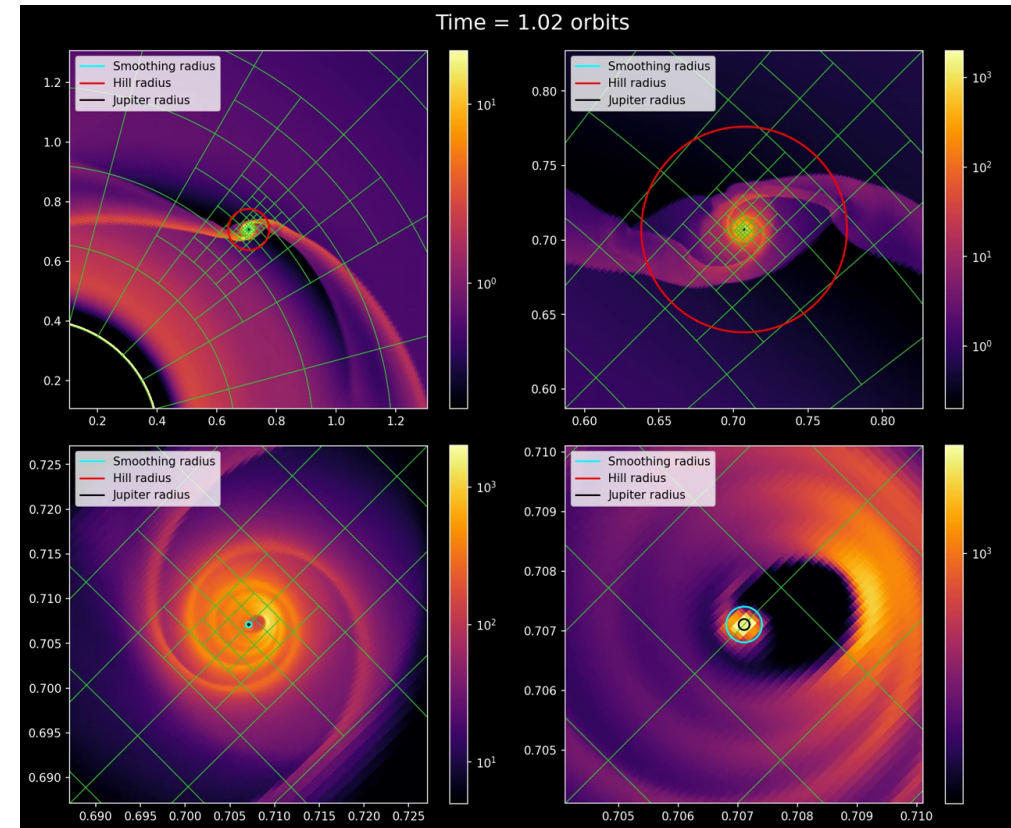


Load balancing, colors are ranks

[Melon Fuksman]

Non-Cartesian geometry AMR

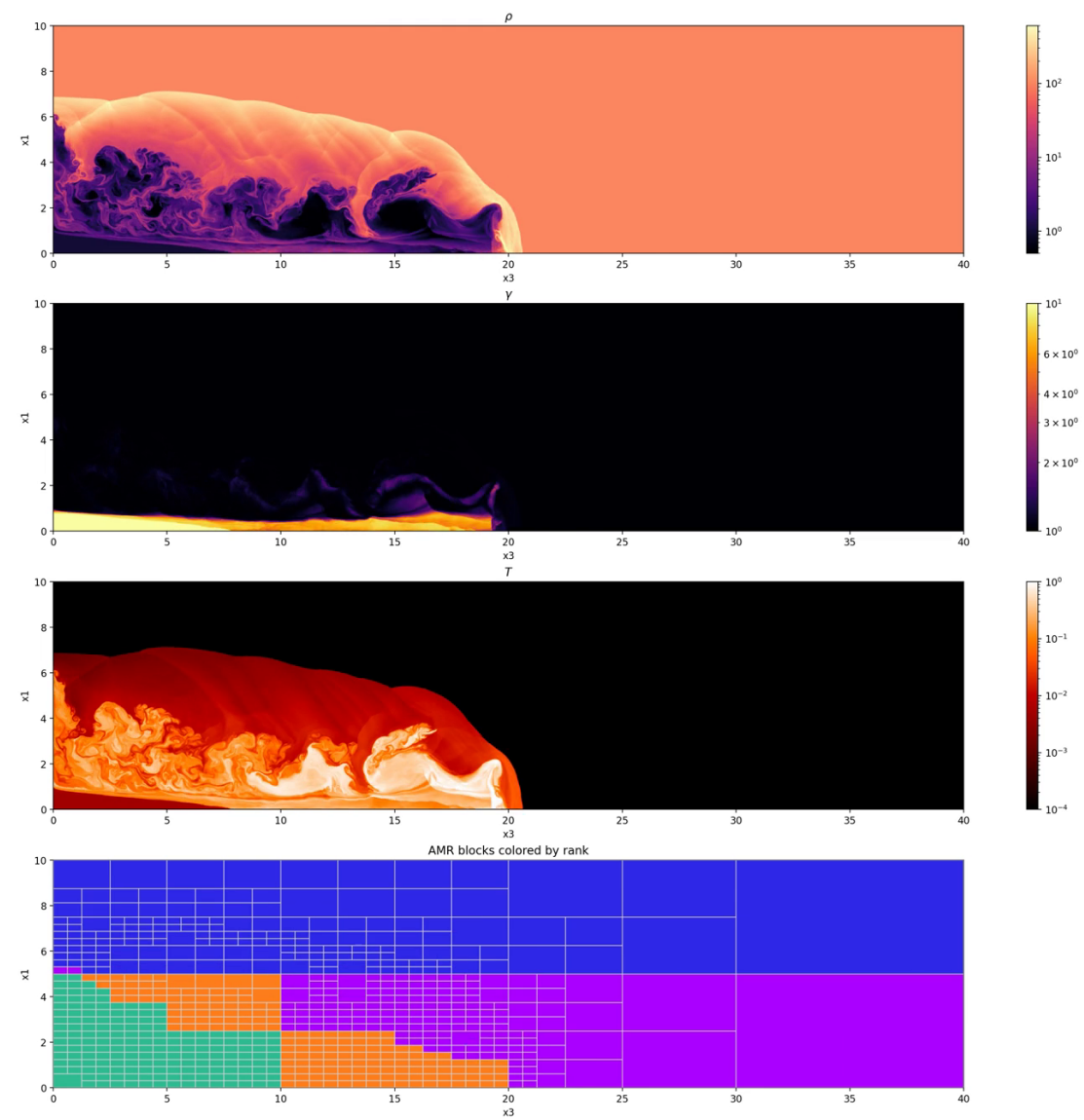
- 10 refinement levels
- similar strategy for oct-tree as cartesian
- Effective resolution: 32768×49152



[Melon Fuksman]

Planet-disk 2d cylindrical with AMR

Relativistic AMR — the jet



relativistic jet with a strong shock

[Melon Fuksman]

Solenoidal constraint

- no magnetic monopoles!
- $\partial_t(\nabla \cdot \mathbf{B}) = -\nabla \cdot (\nabla \times \mathbf{E}) = 0$

But numerically

- discrete fluxes don't enforce $\nabla \cdot \mathbf{B} = 0$
- also truncation error at every step

Issues

- spurious force $\propto (\nabla \cdot \mathbf{B}) \mathbf{B}$ – plasma pushed along $\mathbf{B}$
- wrong jump conditions, negative pressures, crashes

Solutions

- Source terms: Powell's eight-wave
- Divergence cleaning (GLM)
- Constrained transport – zero by construction

Powell's eight-wave Powell+ (1999)

Add a source $\propto \nabla \cdot \mathbf{B}$ to each affected equation

$$\partial_t(\rho \mathbf{v}) + \dots = -(\nabla \cdot \mathbf{B}) \mathbf{B}$$

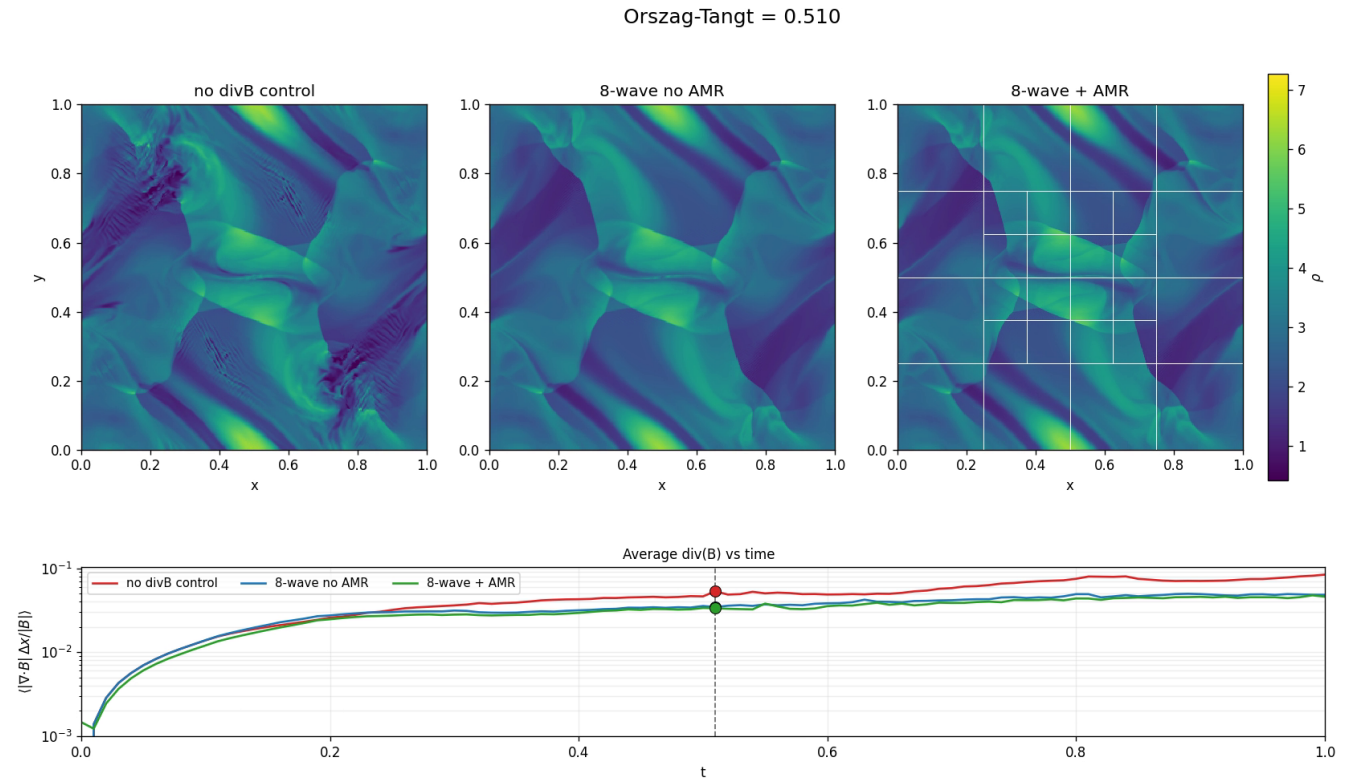
$$\partial_t \mathbf{B} + \dots = -(\nabla \cdot \mathbf{B}) \mathbf{v}$$

$$\partial_t E + \dots = -(\nabla \cdot \mathbf{B}) (\mathbf{v} \cdot \mathbf{B})$$

- Adds an **8th wave** – $\nabla \cdot \mathbf{B}$ is advected with the flow
- local monopole swept off the numerical domain
- simple, robust but **non-conservative**

The Orszag–Tang vortex

- Domain: $[0,1] \times [0,1]$, 256^2 , all boundaries **periodic**
- Uniform Density: $\rho = 25/9$
- Uniform Pressure: $p = 5/3$ (uniform)
- Velocity: $\mathbf{v} = (-\sin 2\pi y, \sin 2\pi x, 0)$
- Magnetic field: $\mathbf{B} = (-\sin 2\pi y, \sin 4\pi x, 0)$
- EOS: ideal, $\gamma = 5/3$
- Scheme: ideal MHD + RK2 + linear reconstruction



Divergence cleaning (GLM)

Couple $\nabla \cdot \mathbf{B}$ to a new scalar field ψ [Dedner+ 2002]:

$$\partial_t \mathbf{B} + \dots = -\nabla \psi$$

$$\partial_t \psi + c_h^2 (\nabla \cdot \mathbf{B}) = -\frac{c_h^2}{c_p^2} \psi$$

- ψ is a generalized Lagrange multiplier sourced by $\nabla \cdot \mathbf{B}$

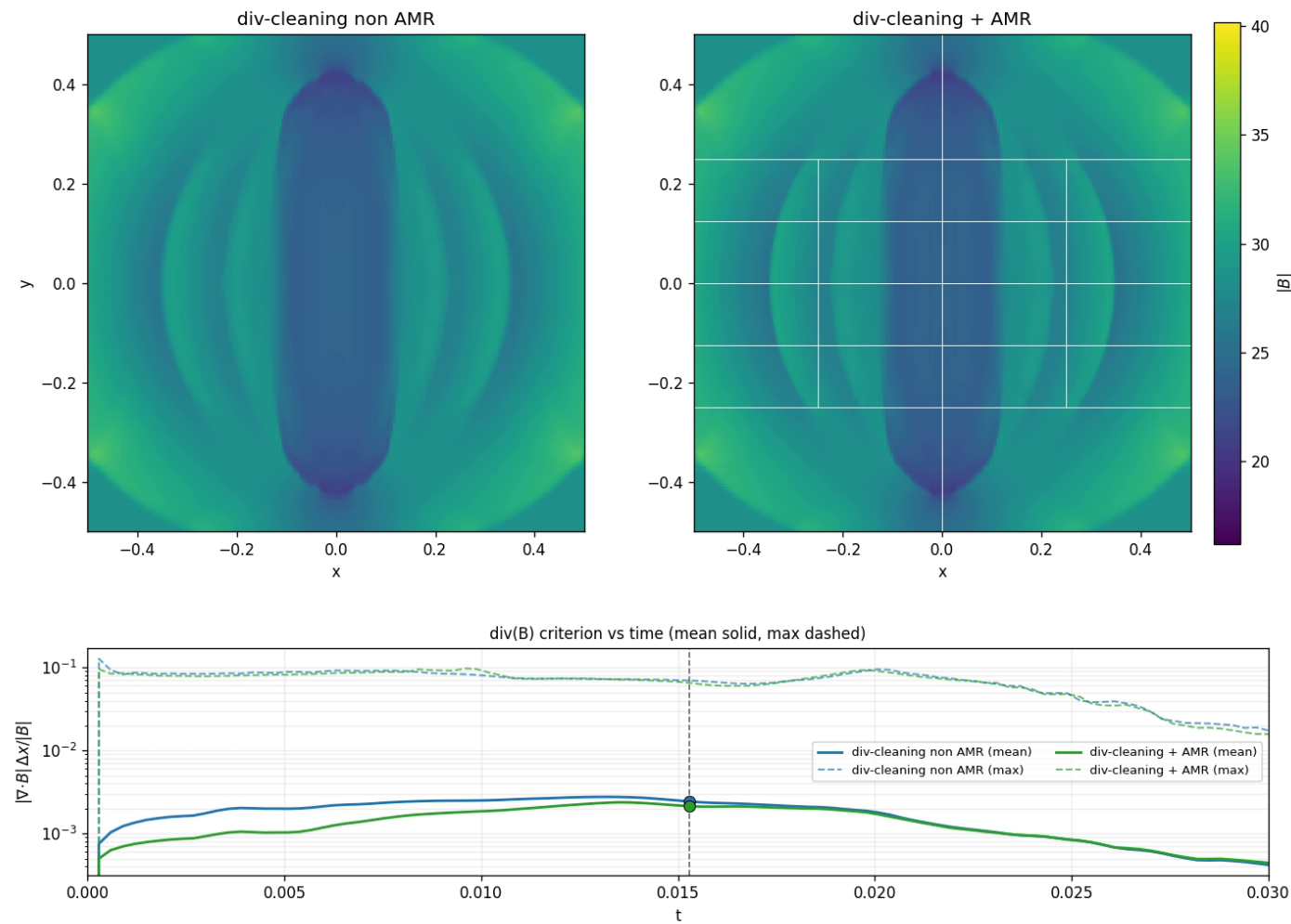
$D \equiv \nabla \cdot \mathbf{B}$ obeys a damped wave equation:

$$\partial_t^2 D = \underset{\substack{\uparrow \\ \text{propagation}}}{c_h^2} \nabla^2 D - \underset{\substack{\uparrow \\ \text{damping}}}{\frac{c_h^2}{c_p^2}} \partial_t D$$

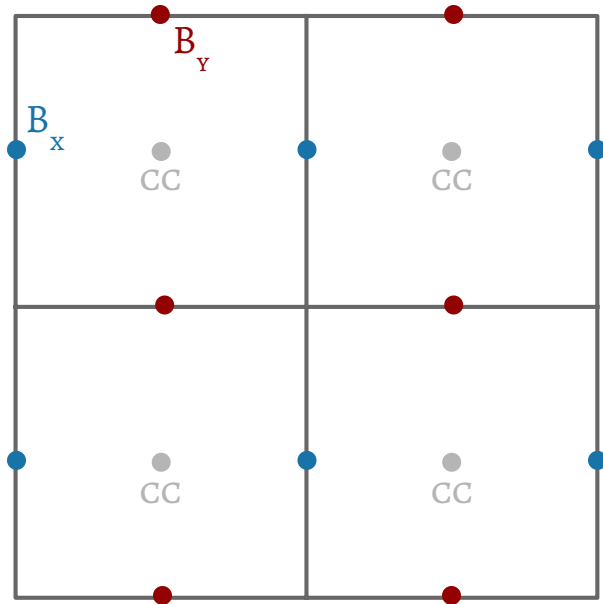
- error is **propagated away** at speed c_h and **damped** on a timescale set by c_p
- conservative and cheap – but only damps; AMR must support this extra equation

Blast wave

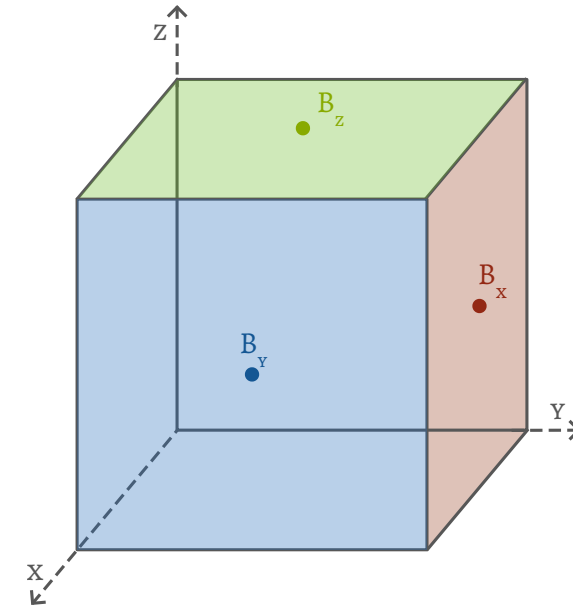
Blast Wave — divergence cleaning $t = 0.0153$



The staggered grid: Magnetic field

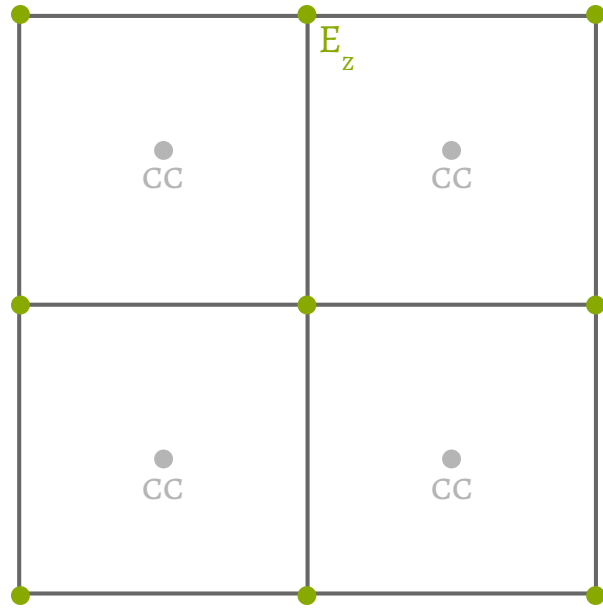


- B_x , B_y at midpoints of x and y faces respectively
- Each face midpoint shared between two neighbouring cells

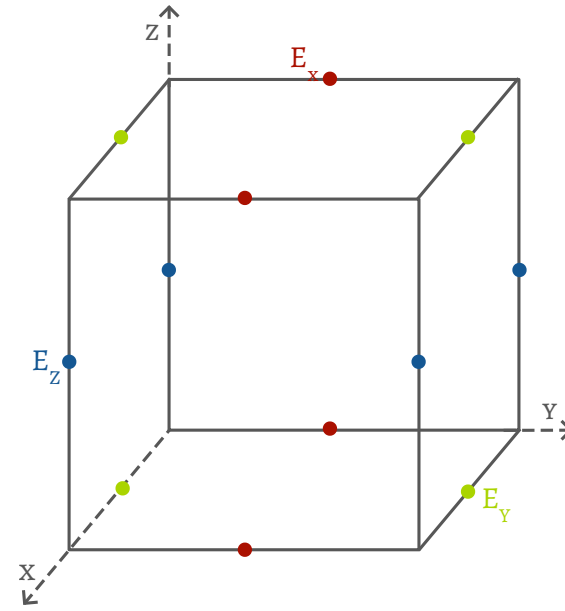


- B_x on yz -faces, B_y on xz -faces, B_z on xy -faces
- Each face value shared between two neighboring cells

The staggered grid: Electric field



- E_z lives at **corners** of grid cells
- Each corner value shared by **4 adjacent cells**



- E_x , E_y , E_z at the midpoint of their own edge – x , y , z edges respectively
- Each edge value shared between four neighboring cells

$\nabla \cdot \mathbf{B} = 0$ preserved to machine precision

For 2D:

$$\frac{\partial B_x^{i+\frac{1}{2},j}}{\partial t} = -\frac{E_z^{i+\frac{1}{2},j+\frac{1}{2}} - E_z^{i+\frac{1}{2},j-\frac{1}{2}}}{\Delta y}, \quad \frac{\partial B_y^{i,j+\frac{1}{2}}}{\partial t} = +\frac{E_z^{i+\frac{1}{2},j+\frac{1}{2}} - E_z^{i-\frac{1}{2},j+\frac{1}{2}}}{\Delta x}$$

Discrete divergence at cell:

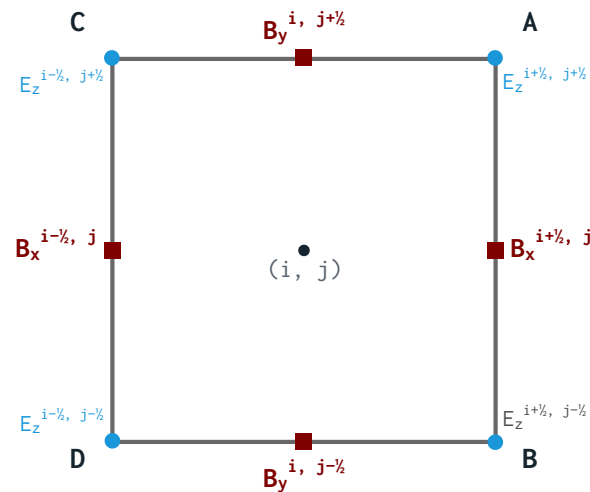
$$(\nabla \cdot \mathbf{B})_{i,j} = \frac{B_x^{i+\frac{1}{2},j} - B_x^{i-\frac{1}{2},j}}{\Delta x} + \frac{B_y^{i,j+\frac{1}{2}} - B_y^{i,j-\frac{1}{2}}}{\Delta y}$$

Four corner EMFs: $A = E_z^{i+\frac{1}{2},j+\frac{1}{2}}$, $B = E_z^{i+\frac{1}{2},j-\frac{1}{2}}$, $C = E_z^{i-\frac{1}{2},j+\frac{1}{2}}$, $D = E_z^{i-\frac{1}{2},j-\frac{1}{2}}$

$$\frac{1}{\Delta x} \left(\dot{B}_x^{i+\frac{1}{2},j} - \dot{B}_x^{i-\frac{1}{2},j} \right) = -\frac{A - B - C + D}{\Delta x \Delta y}, \quad \frac{1}{\Delta y} \left(\dot{B}_y^{i,j+\frac{1}{2}} - \dot{B}_y^{i,j-\frac{1}{2}} \right) = +\frac{A - B - C + D}{\Delta x \Delta y}$$

Hence:

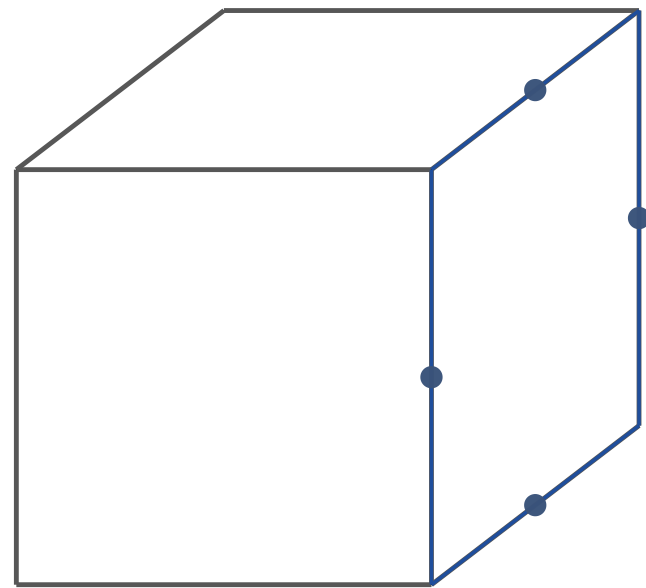
$$\frac{d}{dt}(\nabla \cdot \mathbf{B})_{i,j} = 0$$



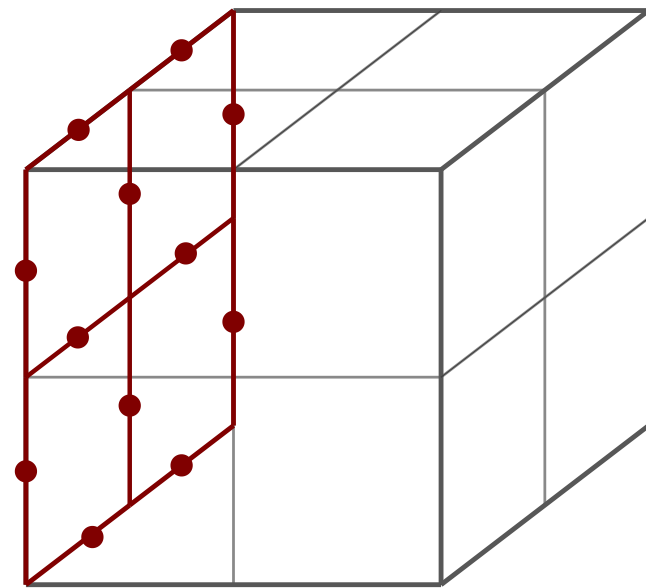
Now add AMR — the real difficulty

- Naïve interpolation of a **face-staggered** field across a coarse-fine boundary **destroys** $\nabla \cdot \mathbf{B} = 0$
- Need **divergence-free prolongation & restriction** for face-allocated fields
- And a **single, consistent EMF** on edges shared between levels **[Stone et. al. 2020]**

EMF correction for face sharers

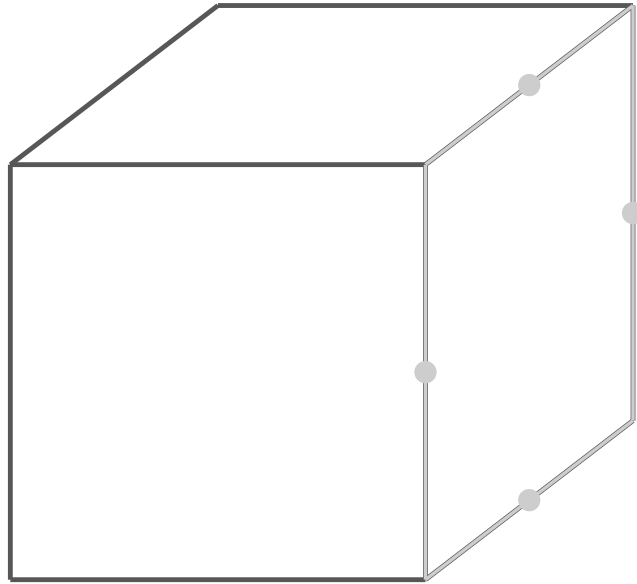


coarse



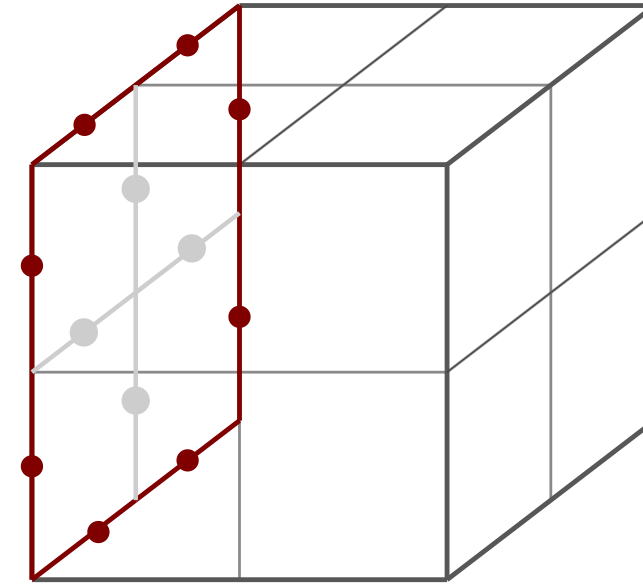
fine

EMF correction for face sharers



coarse

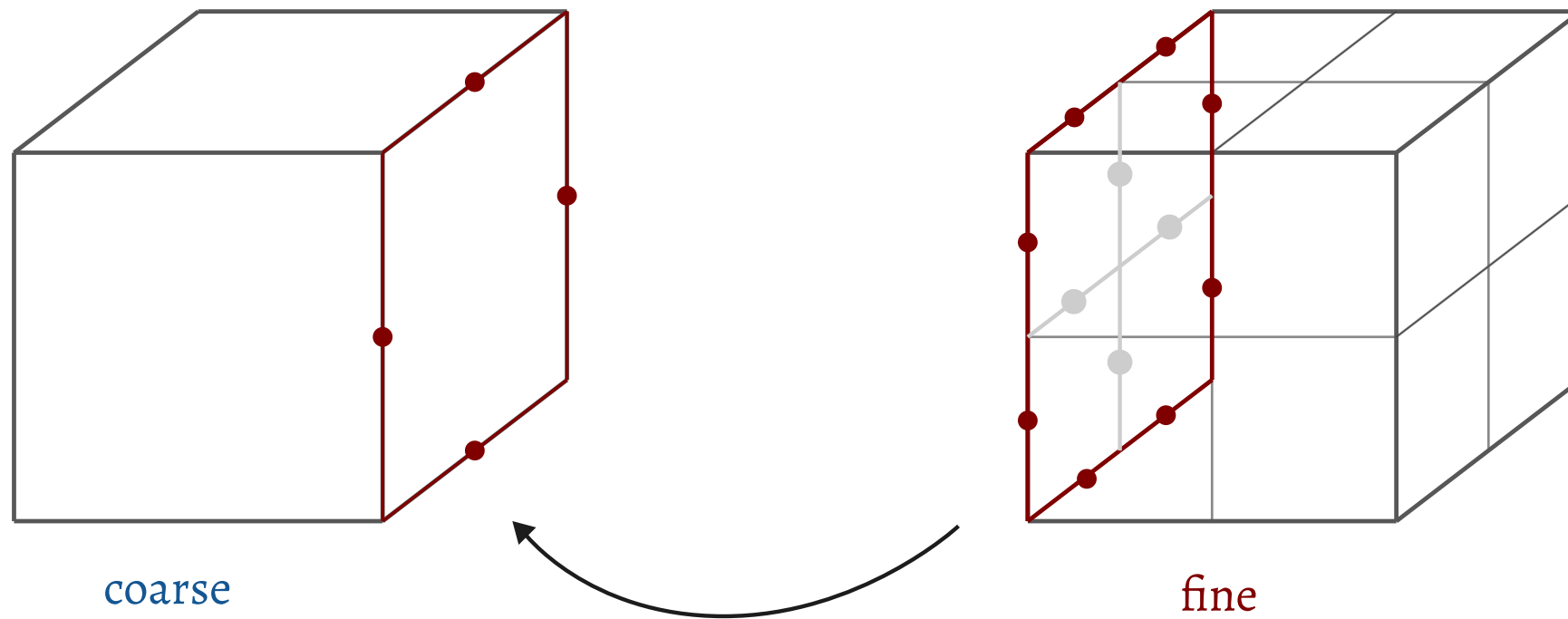
- discard coarse values



fine

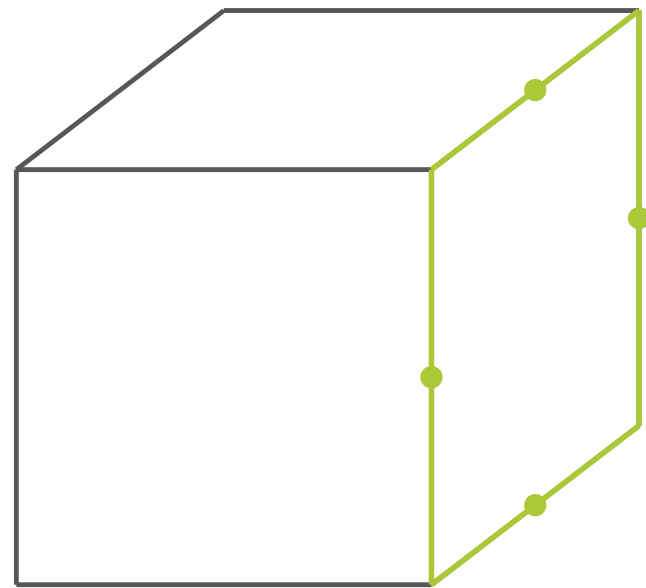
- discard non-overlapping edges

EMF correction for face sharers

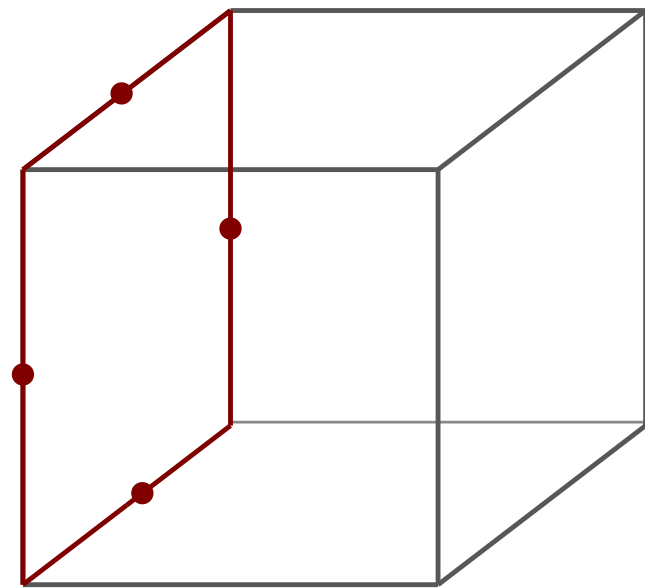


$$\varepsilon_{\text{coarse}} \Delta l_{\text{coarse}} = \sum \varepsilon_{\text{fine}} \Delta l_{\text{fine}}$$

EMF correction for face sharers

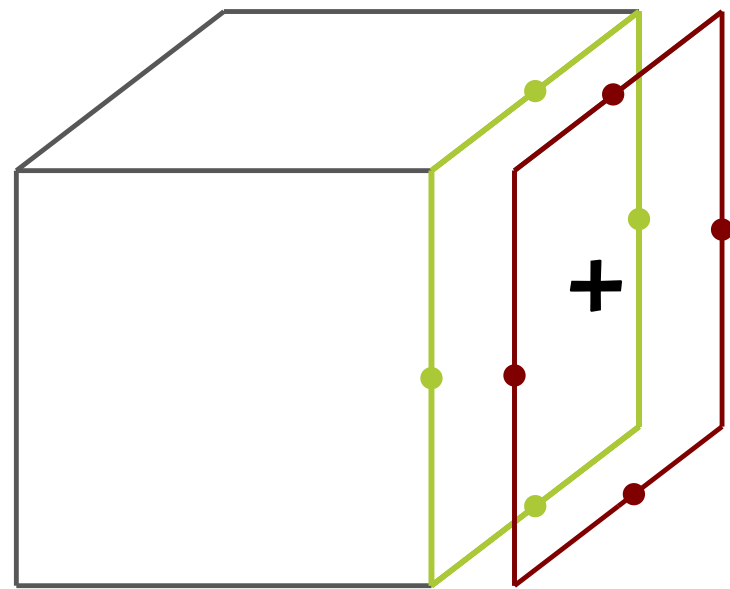


same level

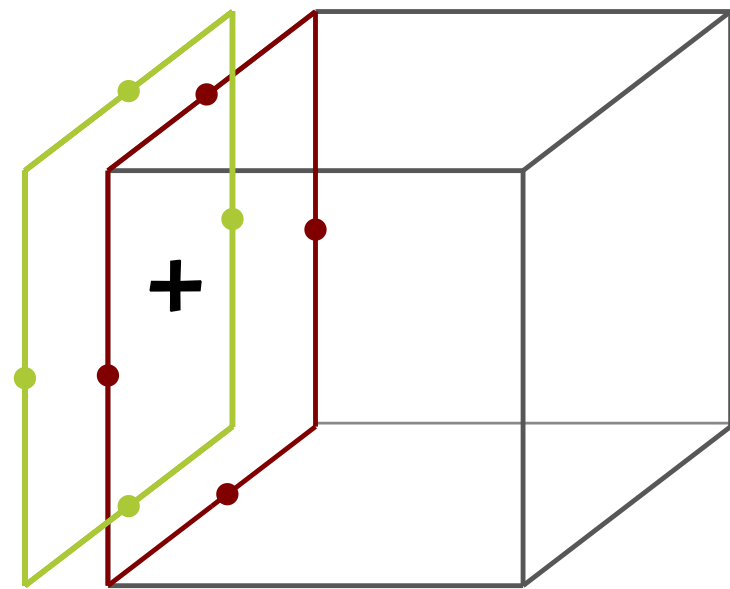


same level

EMF correction for face sharers



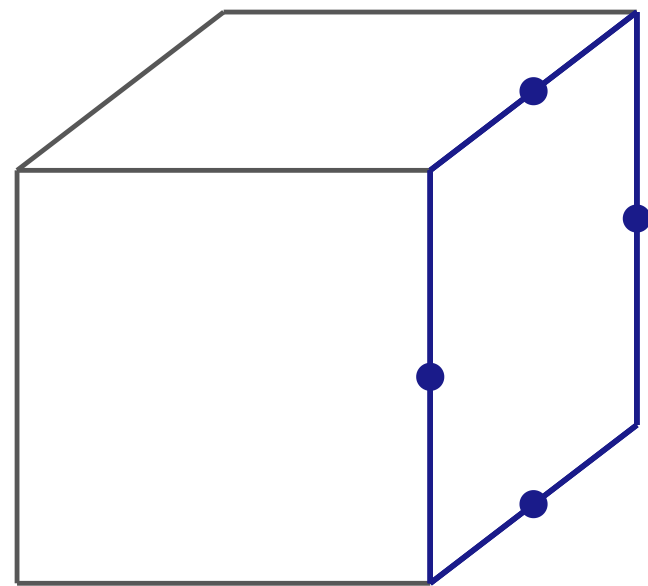
same level



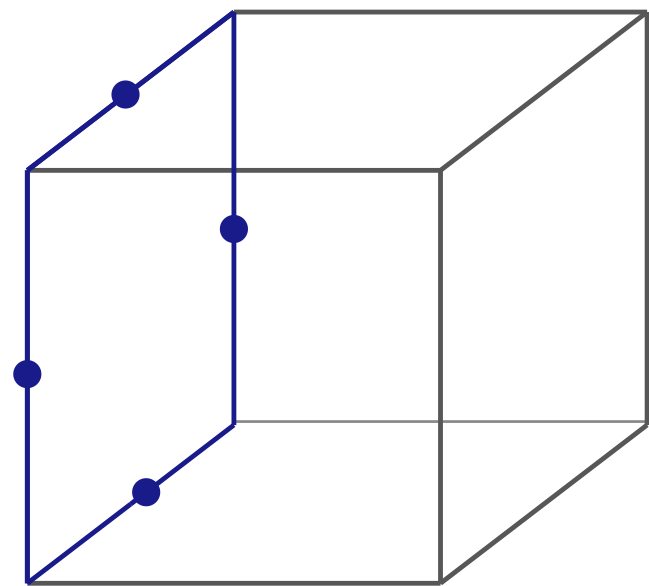
same level

- add the two shared-edge EMFs

EMF correction for face sharers



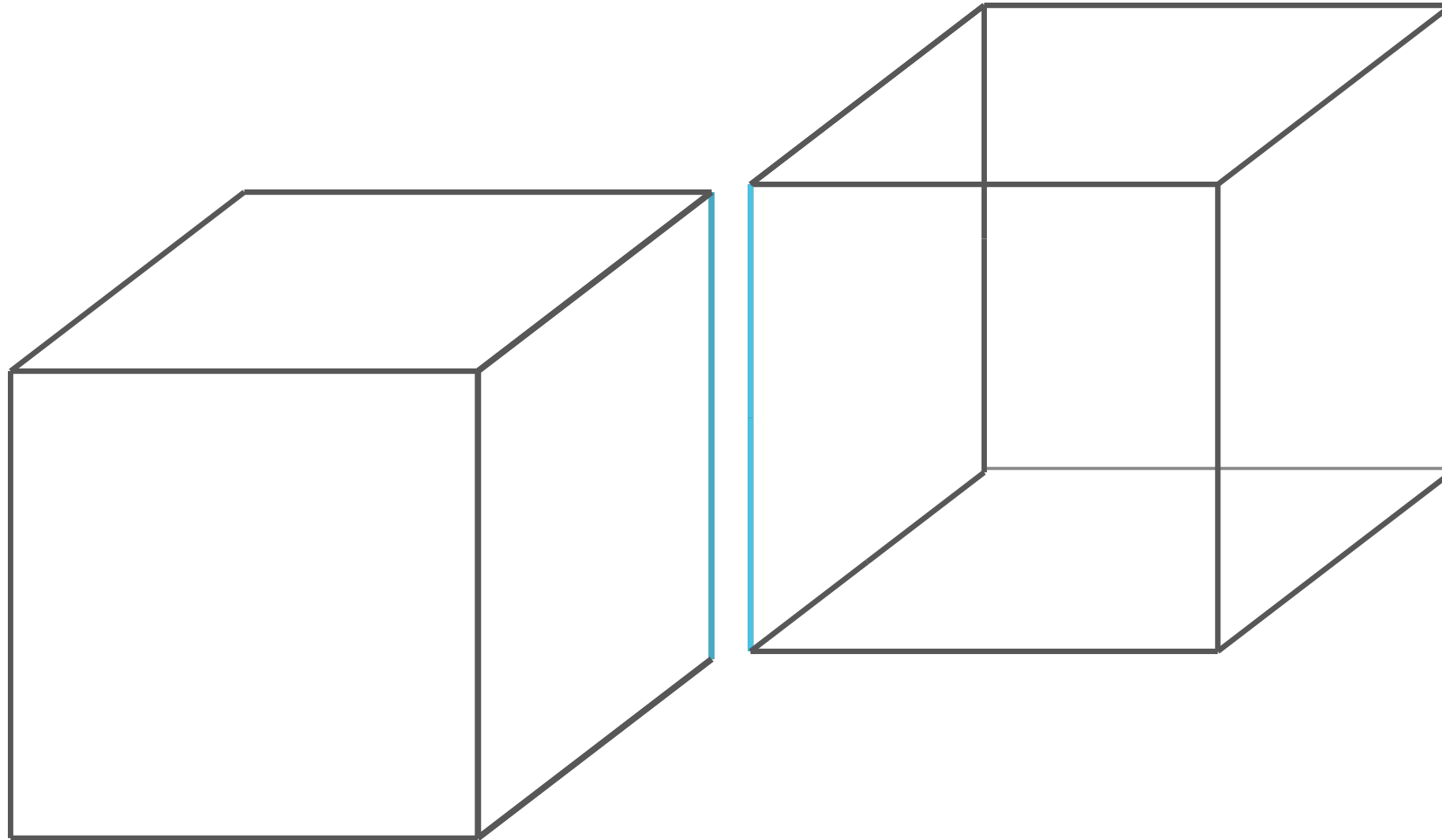
averaged value



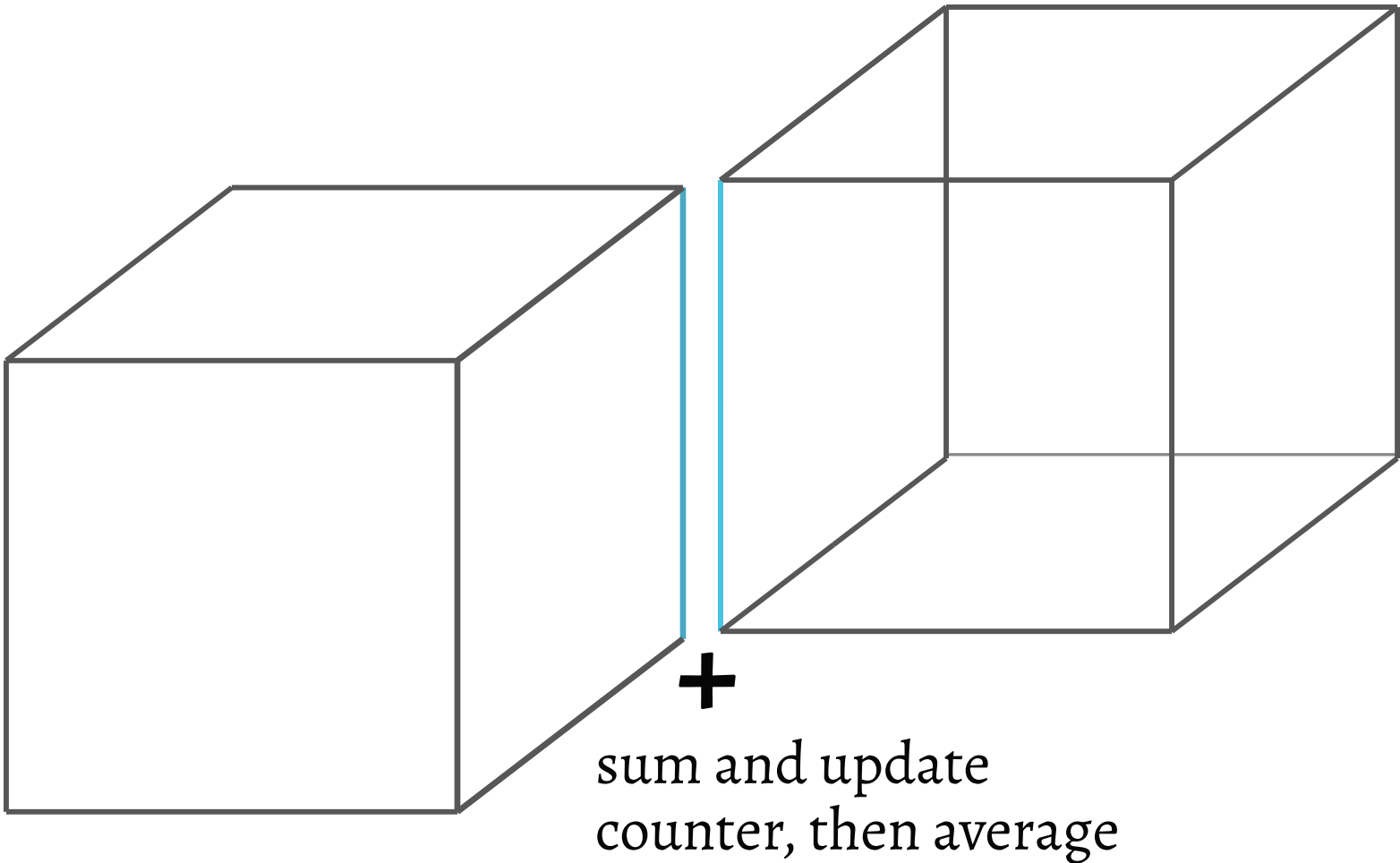
averaged value

$$\varepsilon_L = \varepsilon_R = \frac{1}{2} (\varepsilon_L + \varepsilon_R)$$

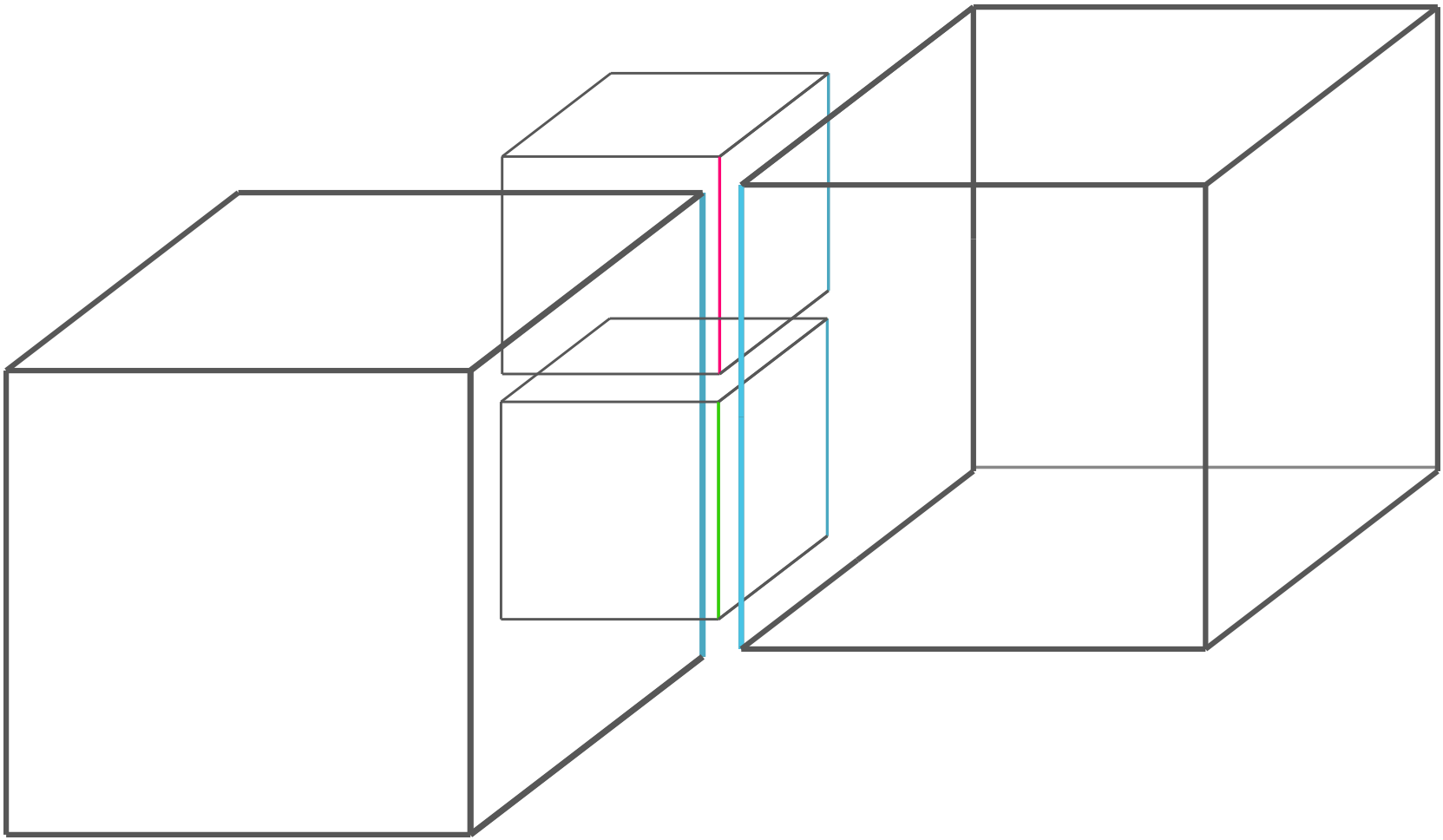
EMF corrections for edge sharers



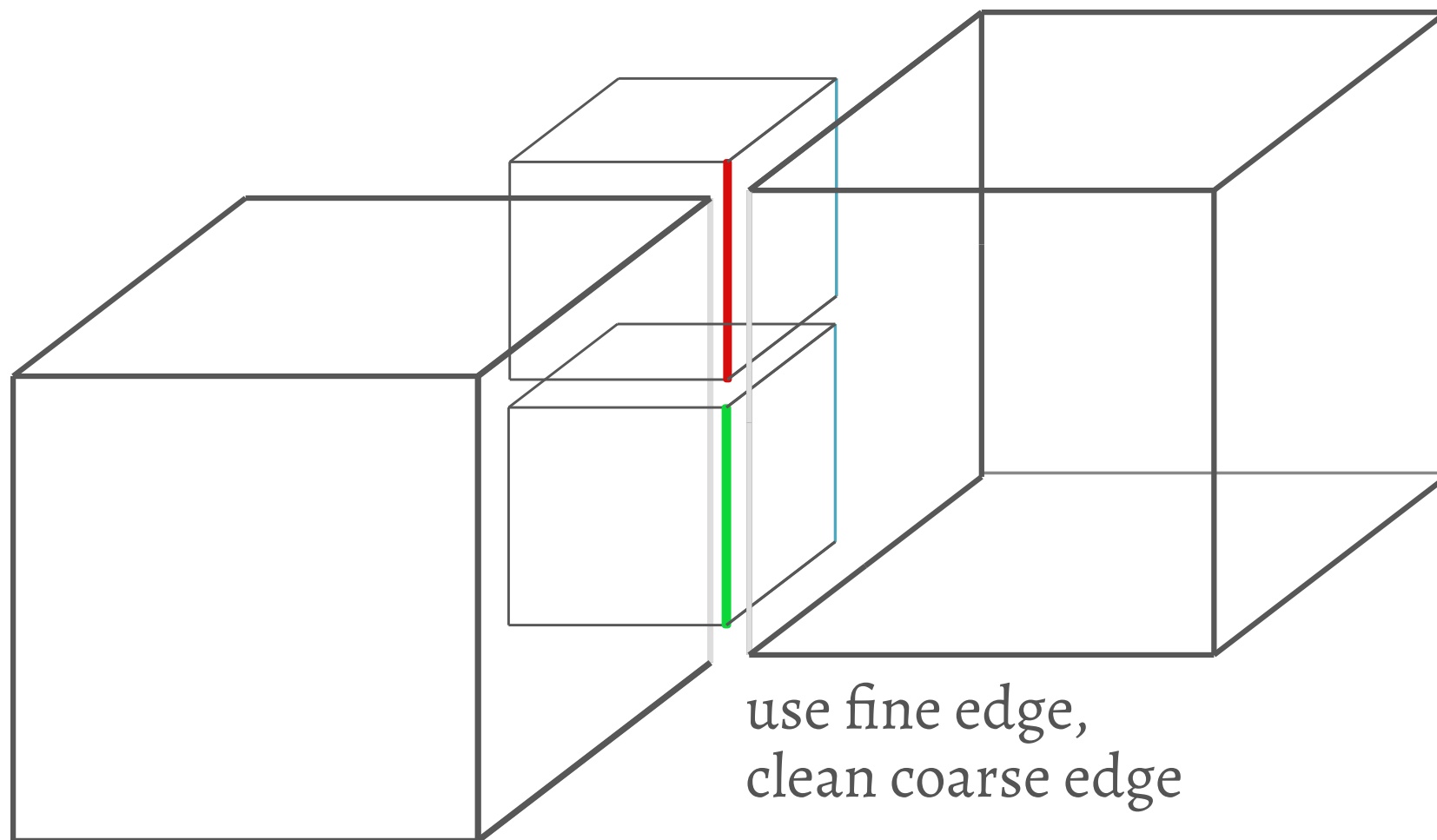
EMF corrections for edge sharers



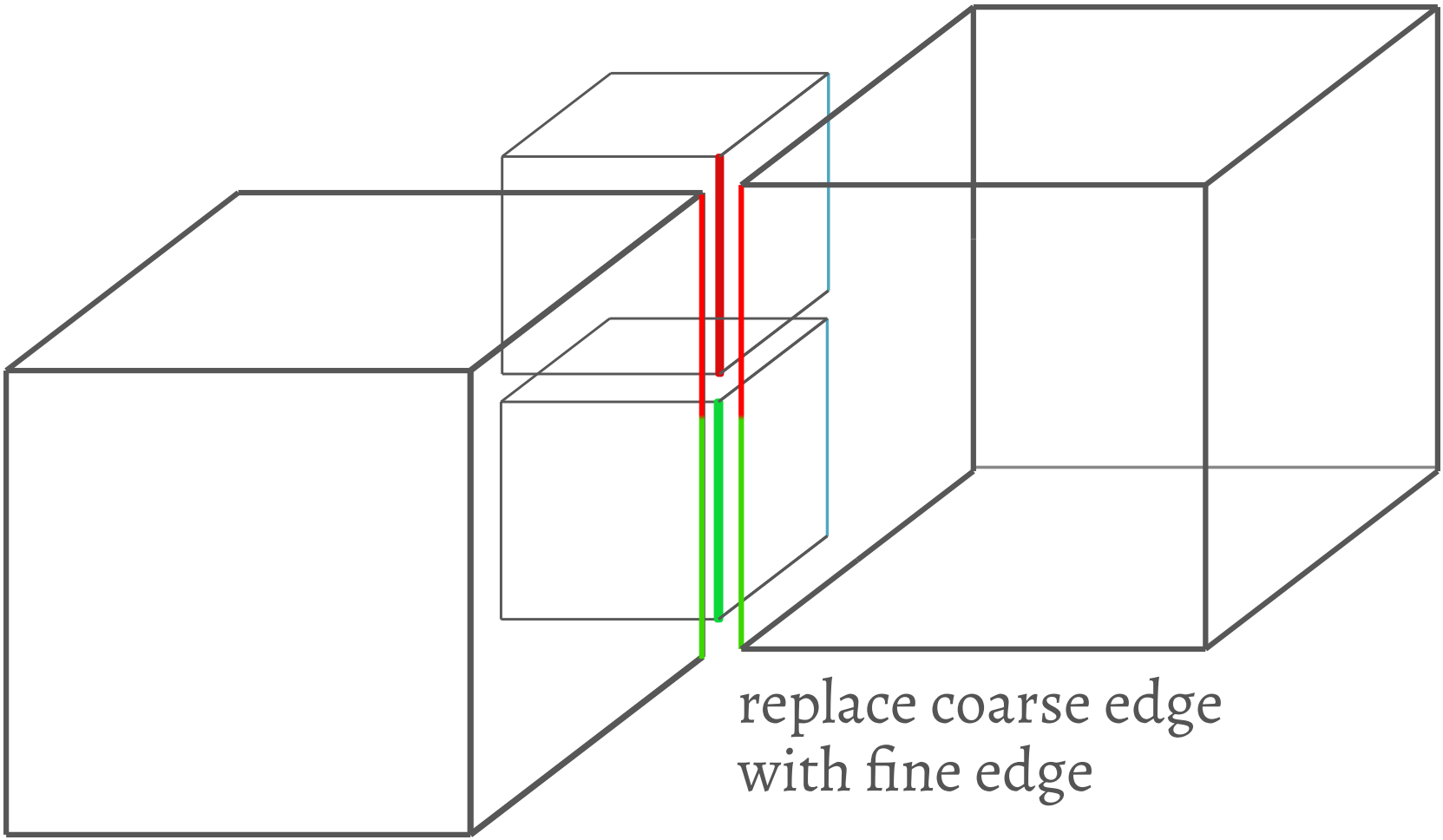
EMF corrections for edge sharers



EMF corrections for edge sharers



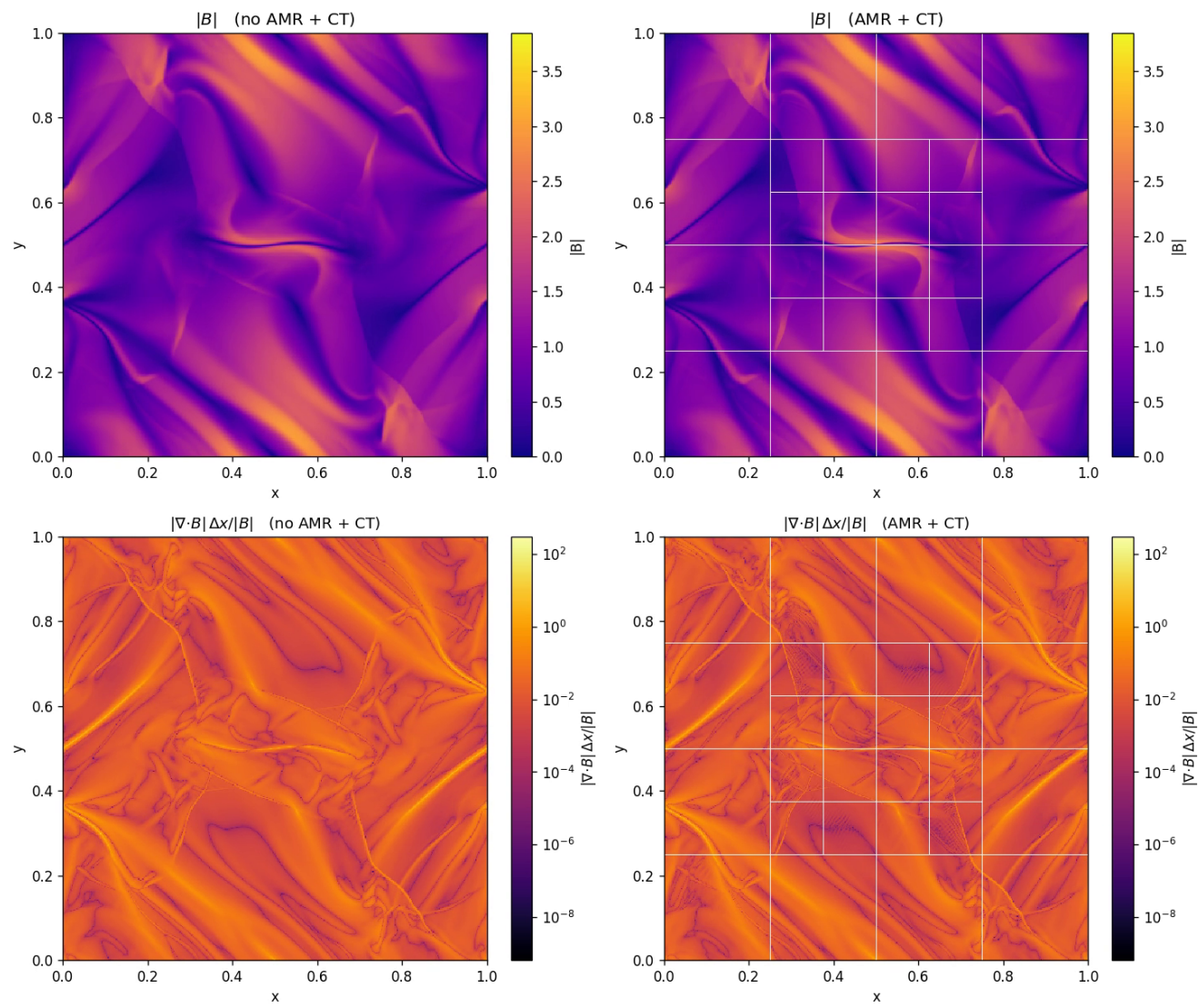
EMF corrections for edge sharers



replace coarse edge
with fine edge

Tests for AMR + CT

Orszag-Tang $t = 0.510$



Thanks

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Questions?