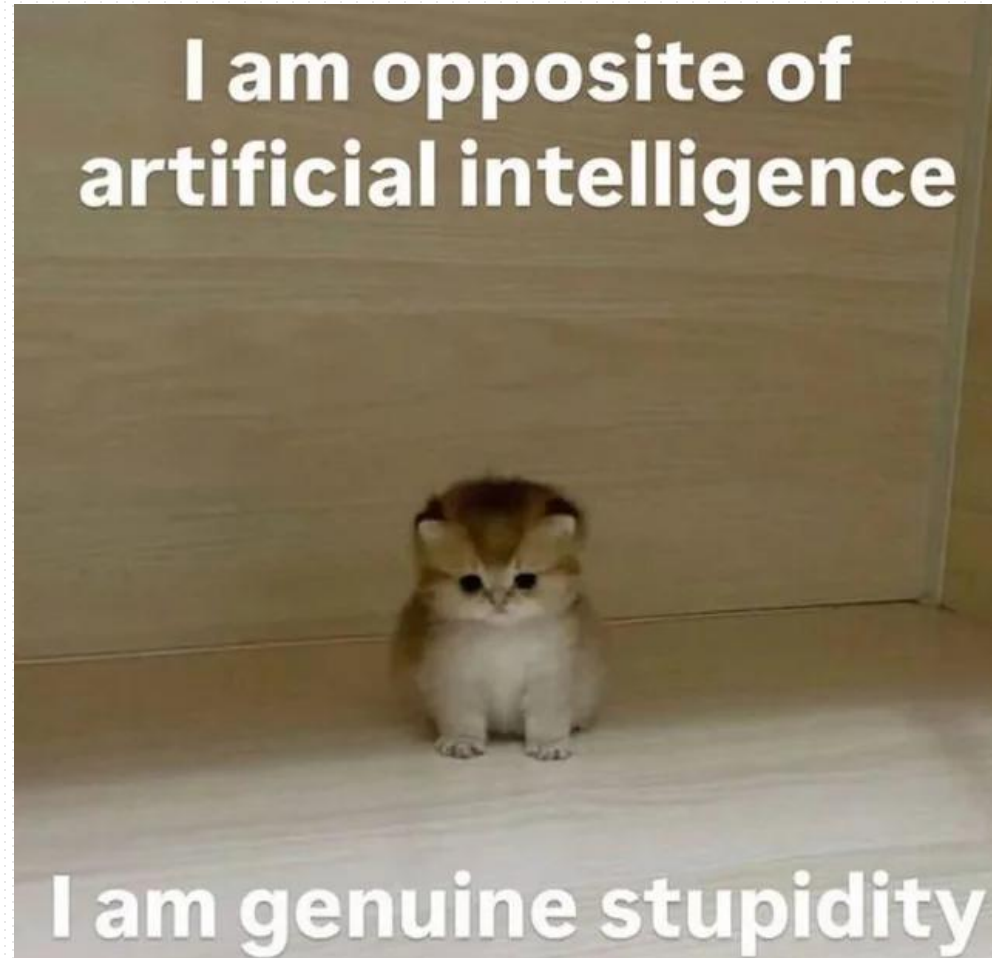


# GRMHD in AthenaK

Jacob Fields

2026 AthenaK Summer School

# A disclaimer

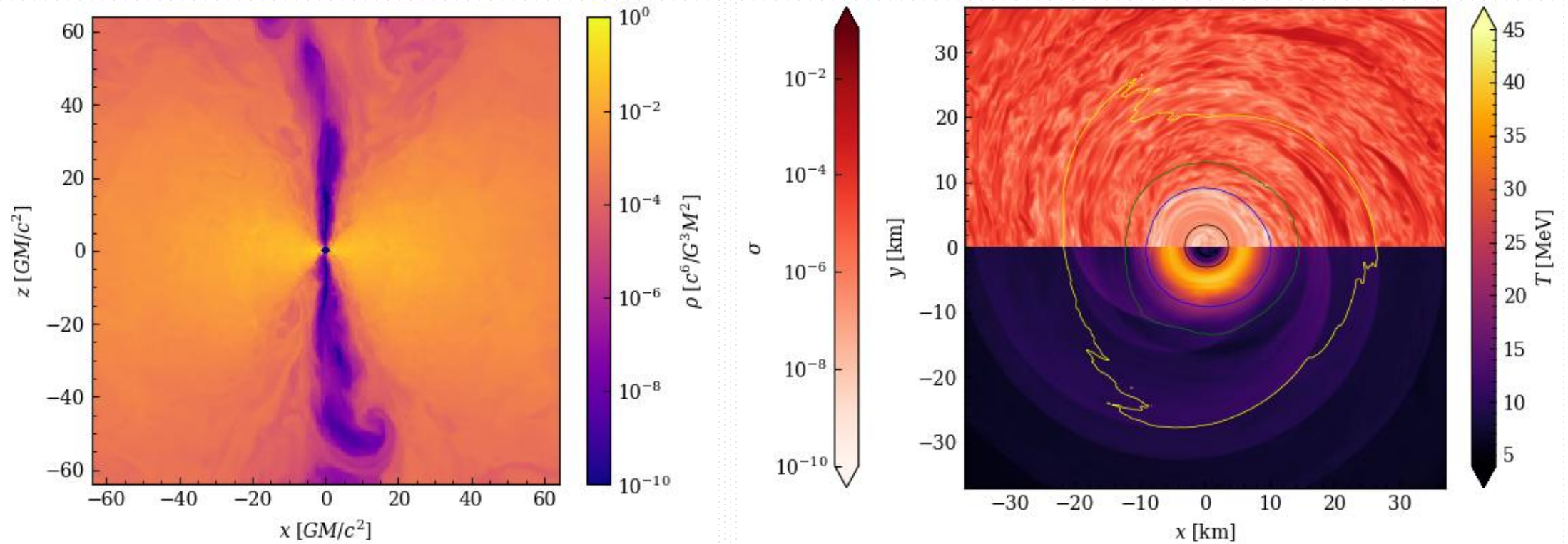


# Challenges with GRMHD

# Where Newtonian MHD breaks down

- Relativistic velocities:  $v \gtrsim 0.1c$
- Gravity is strong:  $\frac{GM}{c^2} \gtrsim 0.1r$
- Relativistic energies:  $\rho\epsilon \sim \rho c^2, B^2 \sim \rho c^2$

# What kinds of problems need GRMHD?



# GRMHD equations

- Conservation of baryon number:

$$\nabla_\mu (\rho u^\mu) = 0$$

- Conservation of stress-energy:

$$\nabla_\mu T^{\mu\nu} = 0,$$

$$T^{\mu\nu} = (\rho h + b^2) u^\mu u^\nu + \left( P + \frac{b^2}{2} \right) g^{\mu\nu} - b^\mu b^\nu$$

- Gauss-Faraday law:

$$\begin{aligned} \nabla_\mu (*F)^{\mu\nu} &= 0, \\ (*F)^{\mu\nu} &= b^\mu u^\nu - b^\nu u^\mu \end{aligned}$$

# Flux-conservative GRMHD equations

- Baryon mass conservation:

$$\partial_t(\sqrt{-g}\rho u^0) + \partial_i(\sqrt{-g}\rho u^i) = 0,$$

- Momentum and energy conservation:

$$\partial_t(\sqrt{-g}T_\mu^t) + \partial_i(\sqrt{-g}T_\mu^i) = \frac{1}{2}\sqrt{-g}T^{\alpha\beta}\partial_\mu g_{\alpha\beta}$$

- Faraday's law:

$$\partial_t(\sqrt{-g}\hat{B}^j) + \partial_i(\sqrt{-g}(b^j u^i - b^i u^j)) = 0,$$
$$\hat{B}^i = (*F)^{it}$$

- Gauss's law:

$$\partial_i(\sqrt{-g}\hat{B}^i) = 0$$

## 3+1 form

$$\begin{aligned}
 \partial_t \tilde{D} + \partial_i (\alpha \tilde{D} \hat{v}^i) &= 0, \\
 \partial_t \tilde{S}_j + \partial_i \left( \tilde{S}_j \hat{v}^i - \frac{\alpha b_j \tilde{B}^i}{W} + \alpha \tilde{P}_{\text{tot}} \delta_i^j \right) &= \frac{1}{2} \alpha \tilde{S}^{ik} \partial_j \gamma_{ik} + \tilde{S}_i \partial_j \beta^i - \tilde{E} \partial_j \alpha, \\
 \partial_t \tilde{\tau} + \partial_i \left( \tilde{\tau} \hat{v}^i - \frac{\alpha^2 b^0 \tilde{B}^i}{W} + \alpha \tilde{P}_{\text{tot}} v^j \right) &= \alpha K_{ij} \tilde{S}^{ij} - \tilde{S}^i \partial_i \alpha, \\
 \partial_t \tilde{B}^j + \partial_i (\alpha [\tilde{B}^j \hat{v}^i - \tilde{B}^i \hat{v}^j]) &= 0,
 \end{aligned}$$

where

$$\begin{aligned}
 \tilde{A} &= \sqrt{\gamma} A, & D &= \rho W, & S_i &= (\rho h W^2 + B^2) v_i - (B^l v_l) B_i, \\
 \tau &= \rho h W^2 + B^2 - P_{\text{tot}} - D, & B^i &= n_\mu (*F)^{\mu i}, & \hat{v}^i &= u^i / W
 \end{aligned}$$



# Solving GRMHD numerically

- Still flux-conservative, can just use FV + CT
  - Reconstruct primitive variables to interfaces
  - Solve Riemann problem at interfaces
  - Compute electric fields at edges
  - Add source terms
  - Integrate with RK solver with a stable timestep
  - Compute primitives from conserved variables
- But some of these steps look a little different

# Some things made simpler by GRMHD

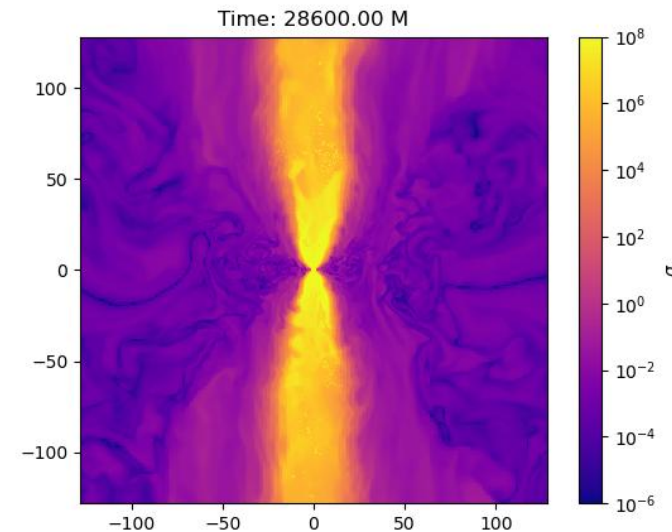
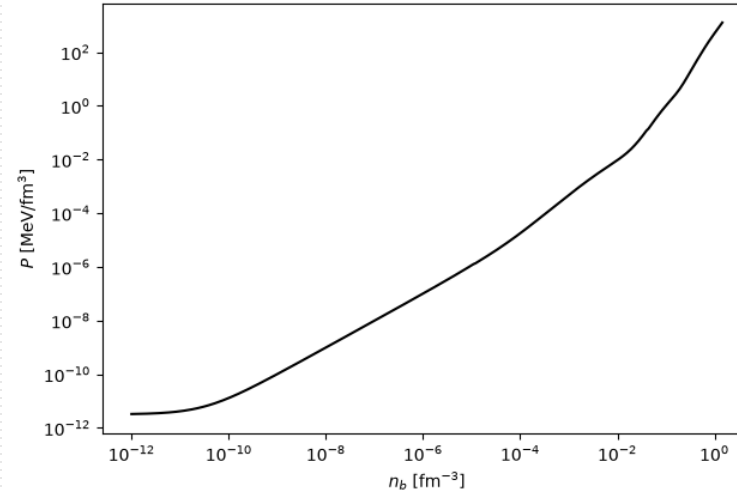
- Gravity is now hyperbolic! (See Manuela and Hengrui's lectures)
  - Arguably easier to solve numerically<sup>1</sup>
- Simpler time step control
  - Maximum characteristic speed, fluid speed never exceeds  $c$
  - Maximum stable timestep  $\Delta t \rightarrow \frac{c\Delta x}{c}$  as  $v \rightarrow c$
- Curvilinear and non-inertial coordinates baked into formalism<sup>2</sup>

# Some things made harder by GRMHD

- Basically everything else.
- Relationship between conserved and primitive variables:
  - $(D, S_i, \tau) \rightarrow (\rho, v^i, P)$  does not have general closed-form solution
  - More constraints on physical state, difficult to detect and repair a priori
- Ill-conditioned at high speeds
  - $W^2 = (1 - v^2)^{-1} \rightarrow \infty$  as  $v \rightarrow 1$
- Characteristic structure more complicated
  - More difficult to construct accurate Riemann solvers
  - Full GRMHD decomposition suitable for numerical implementation unknown until recently (Teukolsky, 2025)\*

# Additional problem-specific challenges

- Binary neutron stars:
  - Nuclear EOS needed
  - Initial data is hard
    - Nonlinear constraints
    - Orbits not quite Keplerian
- Accretion disks
  - $\frac{b^2}{\rho} \gg 1$ , numerically challenging
  - Resolving energetics is challenging



# Some problems have easy solutions

- $v \rightarrow 1$  is ill-conditioned
  - Use  $Wv^i$  wherever possible
- Characteristic structure
  - Use reasonable approximations for wave speeds, more diffusive solvers
- Complex initial data
  - Rely on well-tested and robust external libraries

# Some are harder

- Nuclear EOS
  - Tabulated EOS performs poorly on GPUs
- Conserved-to-primitive problem
  - Need complicated root-solving procedure
  - Need a way to handle unphysical states
- Strongly magnetized regime
  - Code becomes unstable in force-free limit



# Nuclear EOS

- 3D table in  $(n, Y_e, T)$ 
  - $n$  and  $T$  spaced logarithmically
  - Quality of table depends on underlying model, interpolation procedure, etc.
- Three major performance pain points:
  1. Random memory accesses
  2. Often need  $T(n, Y_e, P)$  or  $T(n, Y_e, e)$ ; requires root solve of table
  3. Logs and exponents are expensive

# Floating-point Representation

sign b  
↓  
F  
e

- Floating point
- $m \in [1/2, 1)$  is
- Therefore

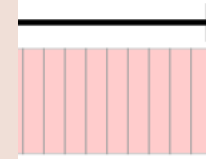


Diagram: Hammond, **JF**,  
Miller, and Barker, *ApJS*  
**277**, 2 (2025)  
([2501.05410](#))

> 0)

xponent

Only thing that needs to be  
calculated

Can be pulled straight from  
the floating-point number



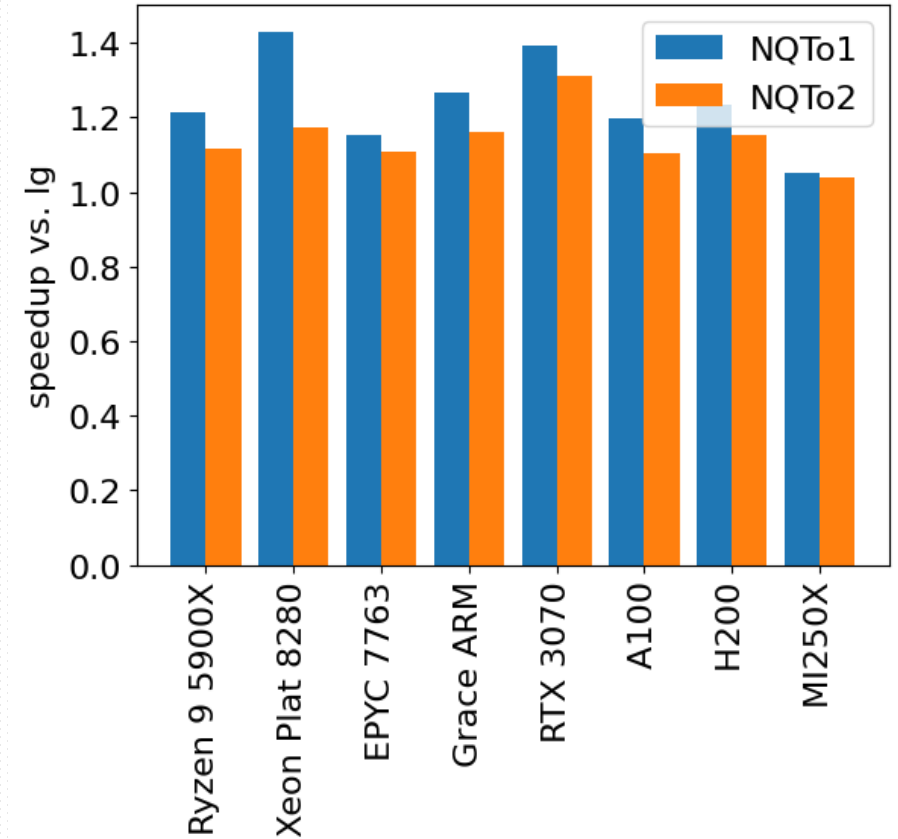
# Not-quite-transcendental functions

- “Log-like” scaling property of table more important than actual scaling
- Can use an approximation for  $\log_2 m$ 
  - Needs to be  $C^1$
  - Needs to be fast
  - Should be exactly invertible

$$\log_2 x \approx f(x) = -\frac{4}{3}(m-2)(m-1) + p$$
$$m = \frac{1}{2} \left( 3 - \sqrt{1 - 3f(m)} \right)$$

# NQT functions (cont.)

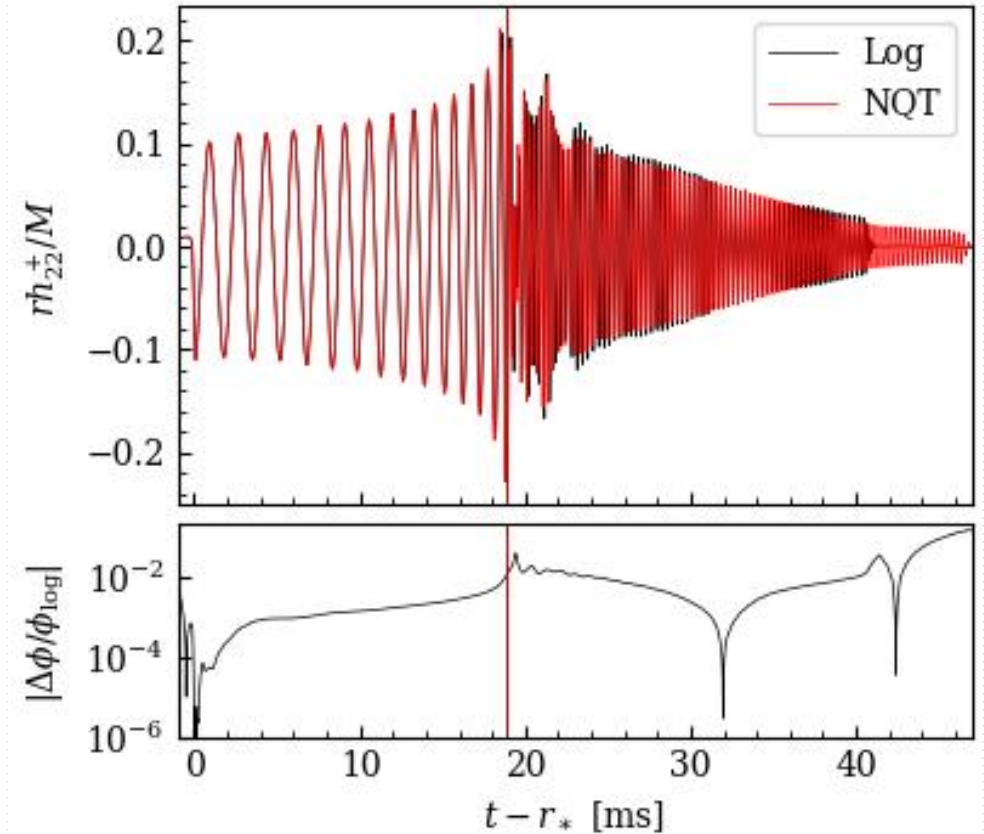
- Implementation in Athenak based on bit hacks for optimal performance
- Table performance improves on all architectures, sometimes up to 30%
- Differs from log-based interpolation by at most a few percent



Hammond, **JF**, Miller, and Barker, *ApJS*  
**277**, 2 (2025) ([2501.05410](#))

# Effects on Waveforms

- NQT-scaled EOS slightly different
  - BNS waveform very sensitive to EOS
- Comparison run at low resolution shows accumulated phase difference is  $<1\%$  during inspiral
- Changing resolution produces a bigger shift in merger time



**JF**, PhD Thesis (2025)

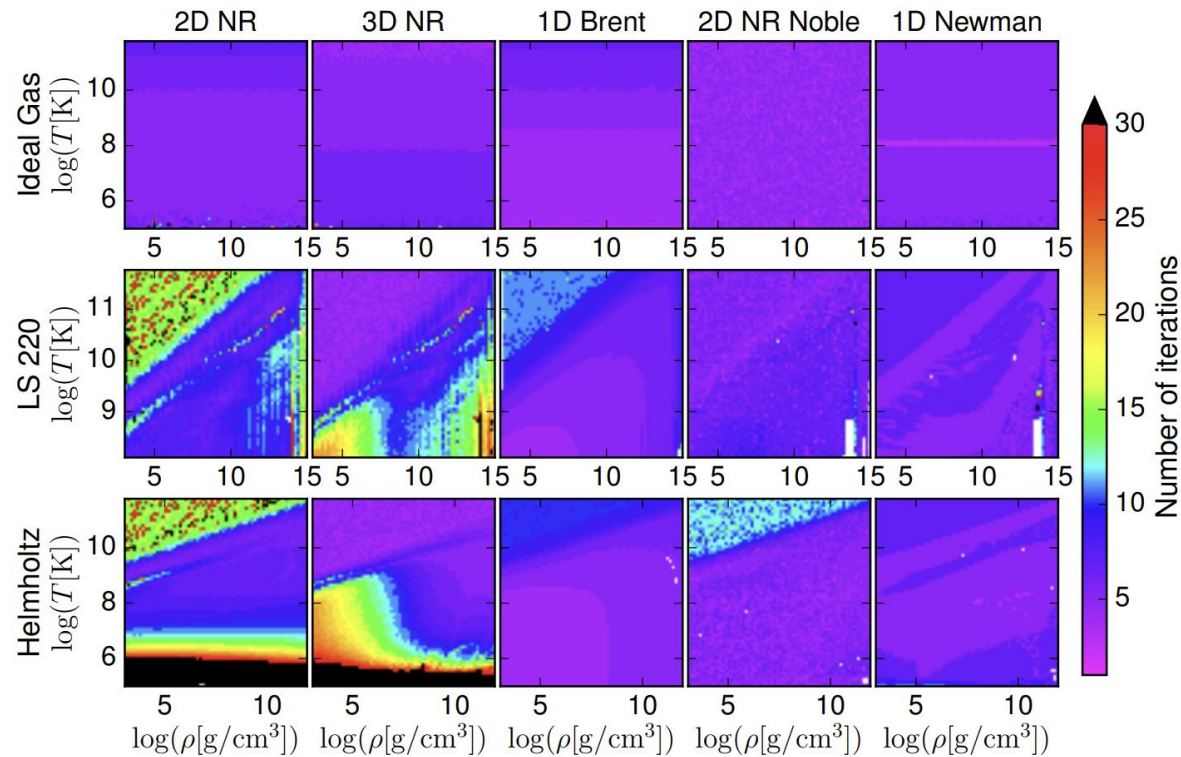
# Conserved-to-primitive inversion

- Need to know  $(\rho, v^i, P)$  to compute fluxes
- Evolved variables are  $(D, S_i, \tau)$  or  $(\rho u^0, T_\mu^t)$
- What makes  $(D, S_i, \tau) \rightarrow (\rho, v^i, P)$  hard?
  - $W^2 = (1 - v^2)^{-1}$  couples directions together
  - $S_i$  depends on EOS via  $h$
- Can generally only be solved iteratively
  - Fast algorithms usually not robust, robust algorithms usually not fast
  - EOS not always smooth => derivative methods not reliable

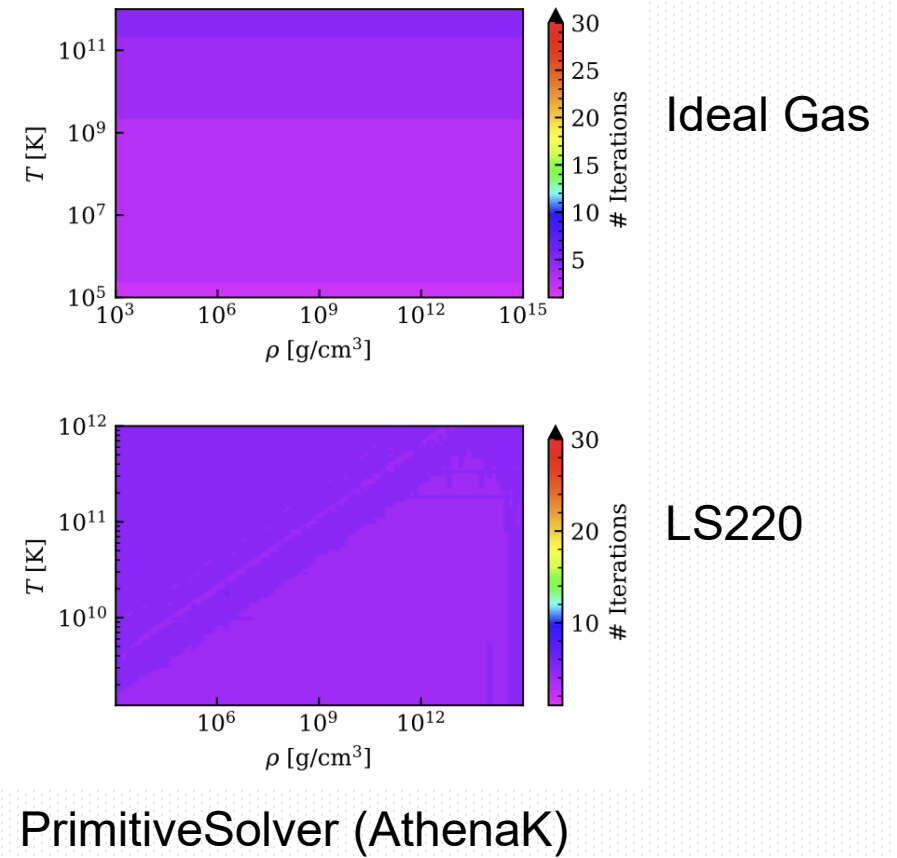
# PrimitiveSolver

- General EOS framework
  - Built-in support for a variety of EOS types
  - Easy to add new EOS
- Implements modified version of C2P inversion from Kastaun *et al.* (2021) (the RePrimAnd algorithm)
  - Start with guess for  $\mu = \frac{1}{hW} \in (0, \frac{1}{h_{\min}})$
  - Use  $\mu$  to find and subtract off B-field component
  - Solve hydro problem for  $\rho$  and  $e$  (or  $\epsilon$ ), refine  $\mu$  guess, rinse and repeat
  - Robust for high  $W$  and  $b^2/\rho$ .

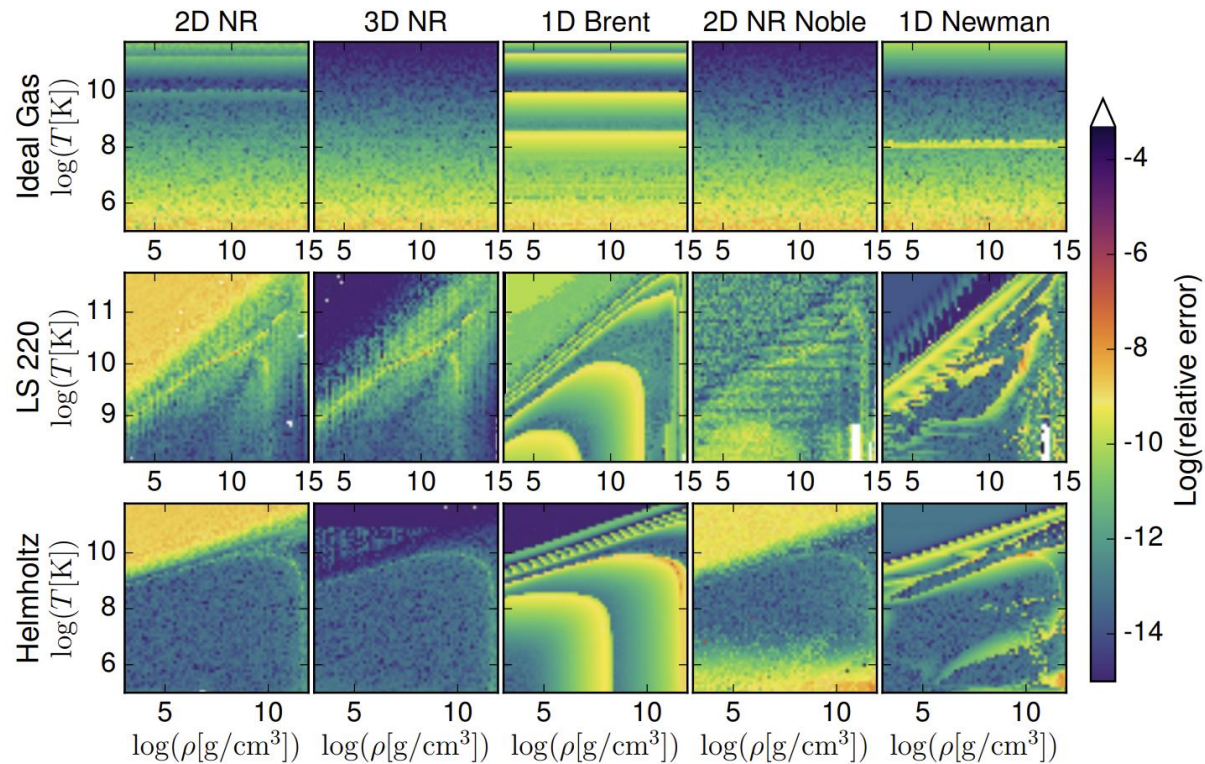
# PrimitiveSolver comparison



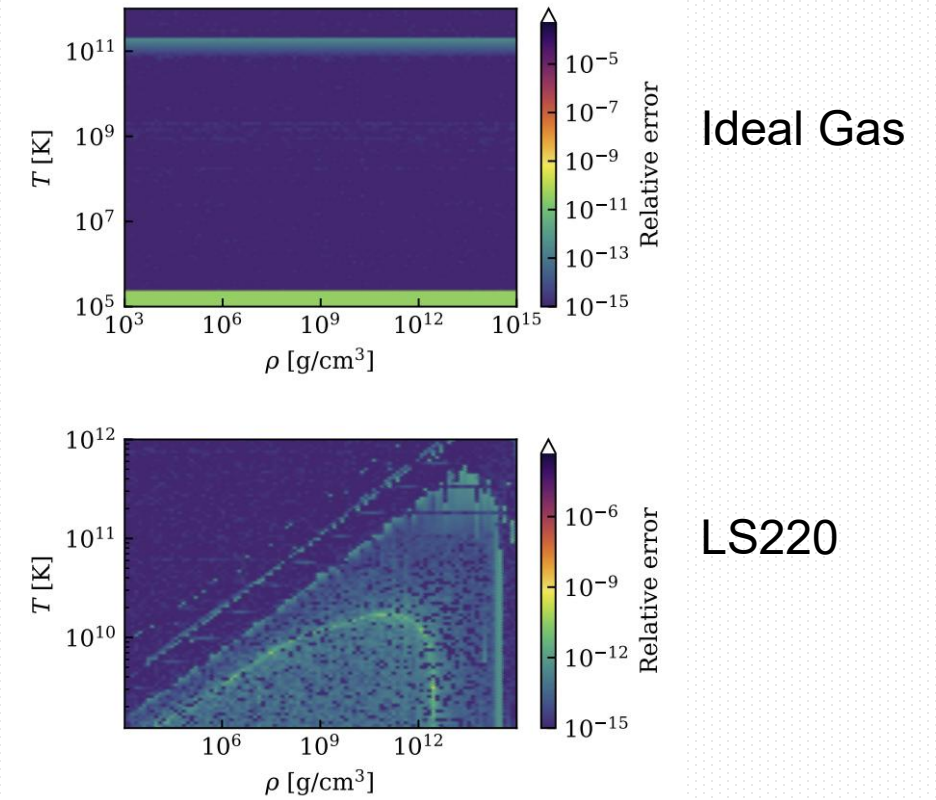
Siegel *et al.* (2018)



# PrimitiveSolver comparison (cont.)



Siegel *et al.* (2018)

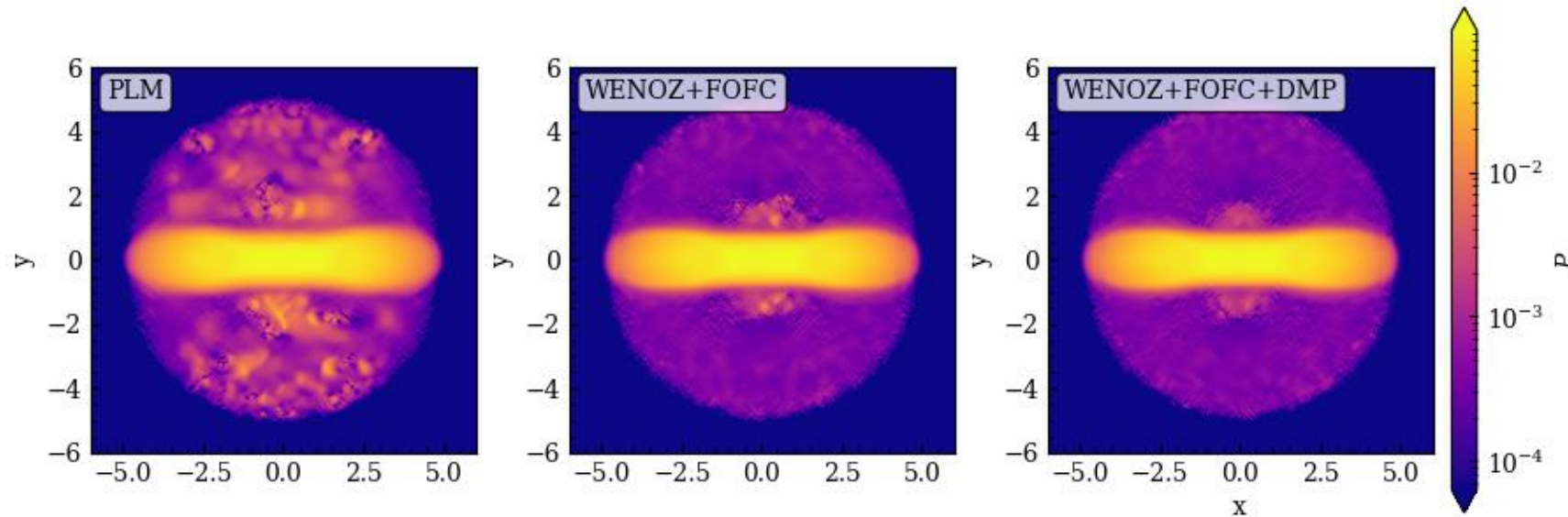


PrimitiveSolver (AthenaK)



# A word on magnetization

- C2P is half the battle: PrimitiveSolver/RePrimAnd is reliable up to extreme magnetizations
- First-order flux correction handles the rest
- Test: strongly magnetized blast wave (Komissarov *et al.*, 1999)





Using GRMHD in AthenaK

# Enabling GRMHD

- Coordinates block:
  - Enable general relativity
  - Set spin (used for CKS spacetime)
  - Enable or disable excision
- Can also use GR solver in flat space; add `minkowski=true` flag to block
- Enables HARM-like solver; only supports CKS/flat space

```
<coord>
general_rel = true      # general relativity
m           = 1.0
a           = 0.0
excise      = false
```

# Example: black hole accretion disk

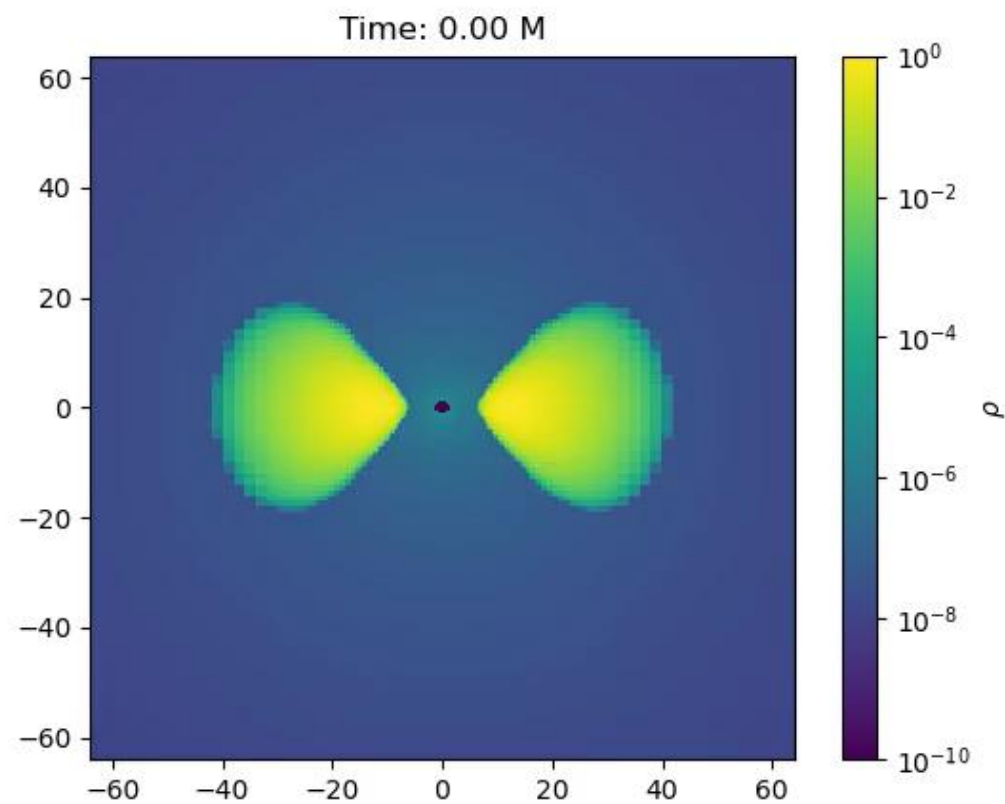
- Compile with `-DPROBLEM=fluids/gr_torus`
- Choose black hole spin and torus configuration (see `inputs/grmhd` for examples)
- Support for both Fishbone-Moncrief and Chakrabarti torii initial conditions

```
<coord>
general_rel = true      # general relativity
a            = 0.9375    # black hole spin a (0 <= a/M < 1)
excise       = true      # excise r_ks <= 1.0
dexcise      = 1.0e-10    # density inside excision
pexcise      = 0.333e-12  # pressure inside excision

<mhd>
eos          = ideal     # EOS type
reconstruct  = ppm4      # spatial reconstruction method
rsolver      = hllc      # Riemann-solver to be used
dfloor       = 1.0e-10    # floor on density rho
pfloor       = 0.333e-12  # floor on gas pressure p_gas
gamma        = 1.3333333333333333 # ratio of specific heats Gamma
fofc         = true      # Enable first order flux correction
gamma_max    = 20.0       # Enable ceiling on Lorentz factor

<problem>
fm_torus     = true      # Fishbone & Moncrief
r_edge       = 6.0       # radius of inner edge of disk
r_peak       = 12.0      # radius of pressure maximum; use l instead if negative
tilt_angle   = 0.0       # angle (deg) to incl disk spin axis relative to BH spin in dir of x
potential_beta_min = 100.0 # ratio of gas to magnetic pressure at maxima (diff locations)
potential_cutoff = 0.2    # amount to subtract from density when calculating potential
potential_rho_pow = 1.0   # dependence of the vector potential on density rho
rho_min      = 1.0e-5     # background on rho given by rho_min ...
rho_pow      = -1.5       # ... * r^rho_pow
pgas_min     = 0.333e-7   # background on p_gas given by pgas_min ...
pgas_pow     = -2.5       # ... * r^pgas_pow
rho_max      = 1.0        # if > 0, rescale rho to have this peak; rescale pres by same factor
l            = 0.0        # const ang. mom. per unit mass u^t u_phi; only used if r_peak < 0
pert_amp     = 2.0e-2     # perturbation amplitude
user_hist    = true       # enroll user-defined history function
```

# Example (cont.)



# Enabling Valencia/dynamical GRMHD

- Add ADM or Z4c block
- Specify EOS and error policy
- Swap out `pfloor` with `tfloor`
- Optional:
  - Flooring threshold
  - Some additional options for FOFC (usually unnecessary)
- Supports arbitrary spacetimes

```
<mhd>
# General options
eos          = ideal      # EOS type
reconstruct  = ppmx       # spatial reconstruction method
rsolver      = hlle       # Riemann solver to be used
dfloor       = 1.0e-10    # floor on density rho
gamma        = 2.0        # ratio of specific heats Gamma
fofc         = true

# DynGRMHD specific
dyn_eos      = ideal      # EOS type
dyn_error    = reset_floor # error policy
tfloor       = 1.0e-8
dthreshold   = 1.00       # Threshold for flooring
enforce_maximum = false

<adm>
```

# Supported EOSs

- Ideal gas

$$e(n, T) = nT/(\Gamma - 1)$$

- Piecewise polytropic EOS:

$$e(n, T) = K_i n^{\Gamma_i} / (\Gamma_i - 1) + nT/(\Gamma_{\text{th}} - 1)$$

- Hybrid tabulated EOS:

$$e(n, T) = e_c(n) + nT/(\Gamma_{\text{th}} - 1)$$

- Microphysical tabulated EOS:

$$e(n, Y_e, T)$$

- Very easy to add more via PrimitiveSolver

# Tabulated EOS support

- [PyCompOSE](#) format; can use any EOS in the [CompOSE](#) database\*
- Enable “not-quite-transcendental” functions for better performance (`mhd/use_NQT=true`)
- Set `mhd/eos_policy=compose`, `mhd/nscalars=1`, and `mhd/table=<table_name>`
- Many pgens (TOV, BNS, etc.) also require 1D table slice for initial data

\*Your mileage may vary

# BNS/BHNS initial data

- Jim's comment that AthenaK requires no external libraries comes with a caveat: all BNS runs require external initial data libraries
  - LORENE: classic pseudospectral code, only reliable for quasicircular BNS
  - FUKA: modern code with support for BBH, BHNS, BNS, arbitrary spins, etc.
  - Elliptica: another modern code with support for BBH, BHNS, BNS, arbitrary spins, etc.
  - sgrid: another BNS library, more modern than LORENE but with fewer features than Elliptica or FUKA



# Example: LORENE BNS

- Compile with `-DPROBLEM=dyn_grmhd/lorene/lorene_bns`
- Provide path to initial data
- Provide 1D EOS information if needed
- Set up BNS trackers
- Set up GW extraction

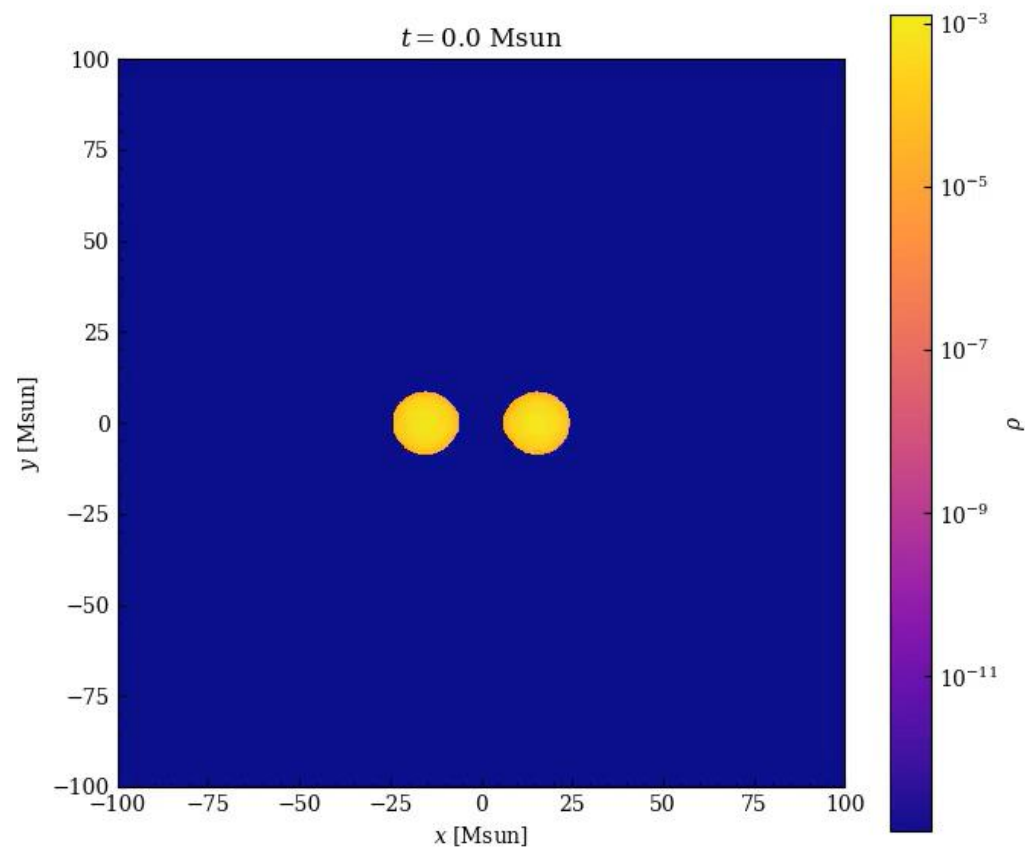
```
# <z4c>
# Gravitational wave extraction
nrad_wave_extraction = 3
extraction_radius_1 = 200
extraction_radius_2 = 400
extraction_radius_3 = 600
waveform_dt = 2.5

# Neutron star tracker.
nco = 2
co_0_type = NS
co_0_x = -15.236509409120682
co_0_radius = 10.0
co_1_type = NS
co_1_x = 15.236509409120682
co_1_radius = 10.0
tracker_mode = walk

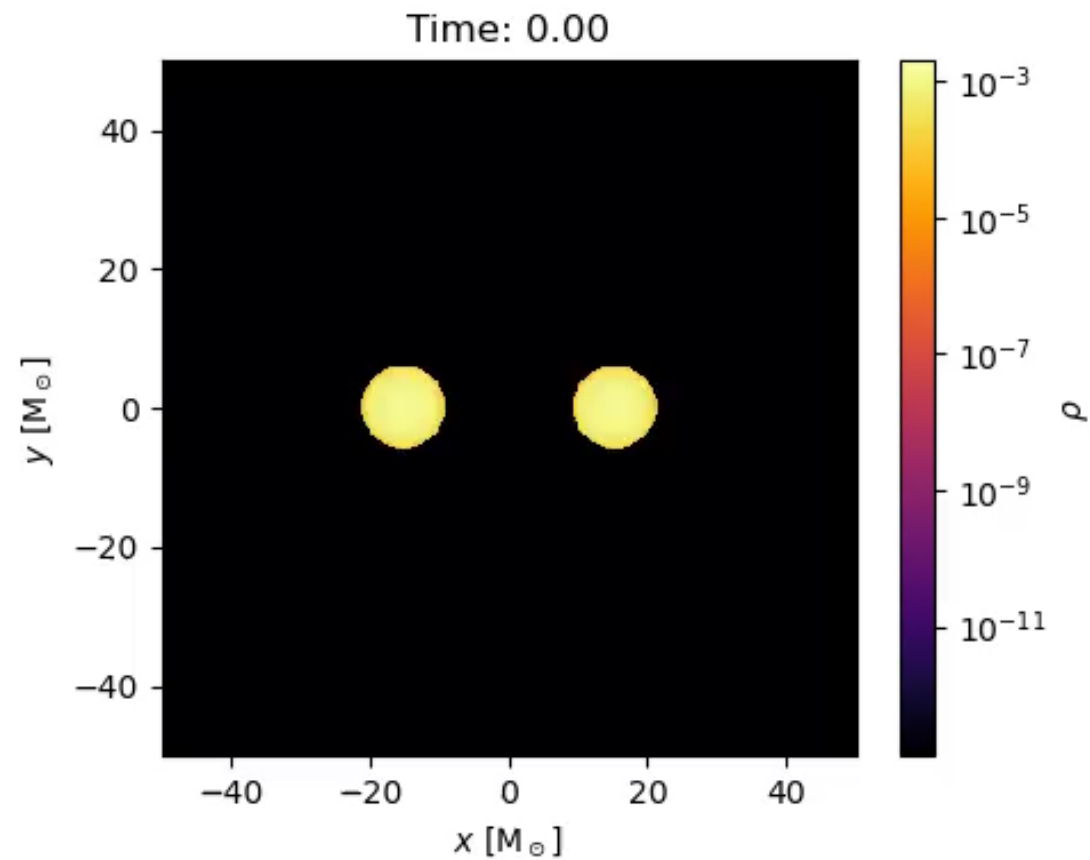
<problem>
# Lorene options
initial_data_file = path_to_lorene.res # Path to LORENE ID
rho_cut = 1e-6 # Cut to atmosphere below this density
b_max = 1e12 # Magnetic field strength in G
r_0_current = 5.0 # Location of current loop for dipole field

# EOS options
kappa = 123.6133 # K for P = K*\rho^\gamma
```

# Example (cont.)



Ideal Gas



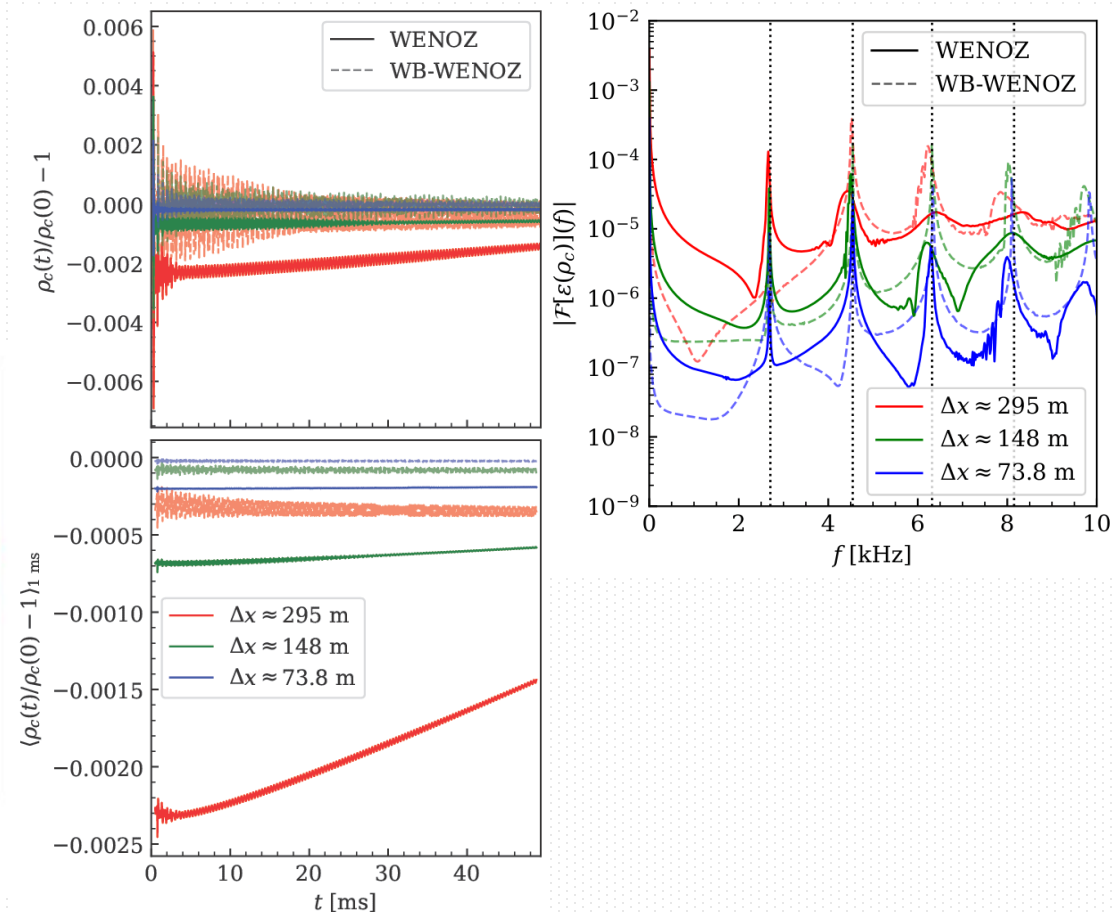
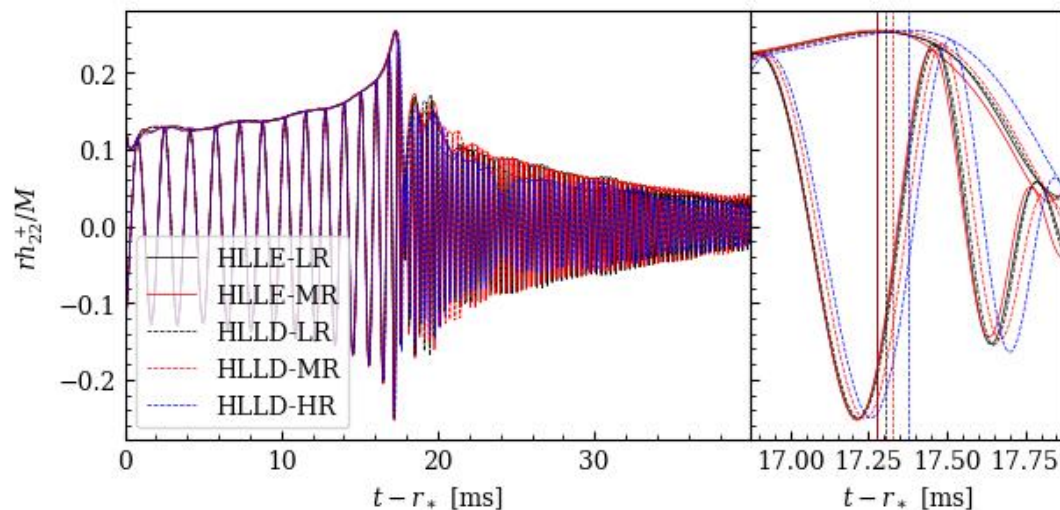
SFHo EOS (tabulated)

# Some options for robustness

- Enable first-order flux correction: `mhd/fofc=true`
- Set a ceiling on  $W$ : `mhd/gamma_max=...`
- Enforce a maximum on  $B^2/D$  (DynGRMHD only):  
`mhd/max_bsq=...`
- Use FOFC to enforce relaxed maximum principle (DynGRMHD only, not usually needed): `mhd/enforce_maximum=true`,  
`dmp_M=...` ( $M \geq 1$ )

# New developments coming soon

- HLLD for GRMHD
- Well-balanced scheme for hydrostatic systems



# Summary

- AthenaK has two GRMHD solvers
  - HARM-like solver (CKS spacetimes)
  - Valencia/3+1 solver (arbitrary spacetimes)
- State-of-the-art methods for speed and stability
  - Robust C2P algorithm
  - First-order flux correction
  - Optionally read NQT-scaled tabulated EOS